

Emergence of Dirac Electron Pair in Charge Ordered State of Organic Conductor α -(BEDT-TTF)₂I₃

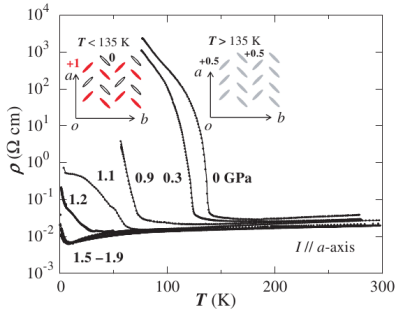
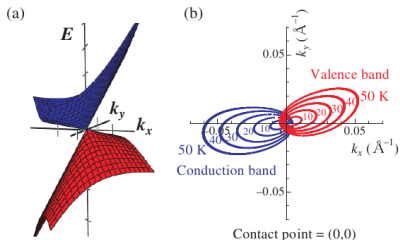
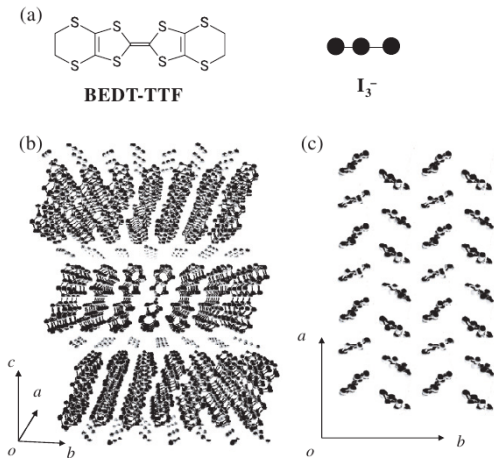
(arXiv:1107.4841)

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The Organic Conductor α -(BEDT-TTF) $_2$ I $_3$



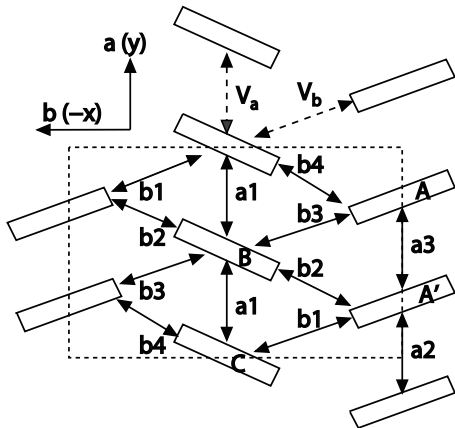
N. Tajima and K. Kajita, *Sci. Tech. Adv. Mater.* **10**, 024308 (2009)

Outline: study of the (V_a, P_a) phase diagram

- Tight-binding model for α -(BEDT-TTF) $_2$ I $_3$ with mean field for e-e interaction
- 2×2 effective Hamiltonian (Luttinger-Kohn representation)
- Berry curvature and Berry phase

The Organic Conductor α -(BEDT-TTF) $_2$ I $_3$

- Insulating stripe Charge Ordered (CO) phase
- Charge Ordered Metallic (COM) phase
- Superconducting state in the presence of charge ordering
- Zero Gap State (ZGS) with a massless Dirac spectrum



N. Tajima and K. Kajita, *Sci. Tech. Adv. Mater.* **10**, 024308 (2009)

Hartree mean-field theory

$$H = \sum_{(i\alpha:j\beta),\sigma} (t_{i\alpha;j\beta} a_{i\alpha\sigma}^\dagger a_{j\beta\sigma} + \text{h.c.}) + \sum_{i\alpha} U a_{i\alpha\uparrow}^\dagger a_{i\alpha\downarrow}^\dagger a_{i\alpha\downarrow} a_{i\alpha\uparrow} \\ + \sum_{(i\alpha:j\beta),\sigma,\sigma'} V_{\alpha\beta} a_{i\alpha\sigma}^\dagger a_{j\beta\sigma'}^\dagger a_{j\beta\sigma'} a_{i\alpha\sigma},$$



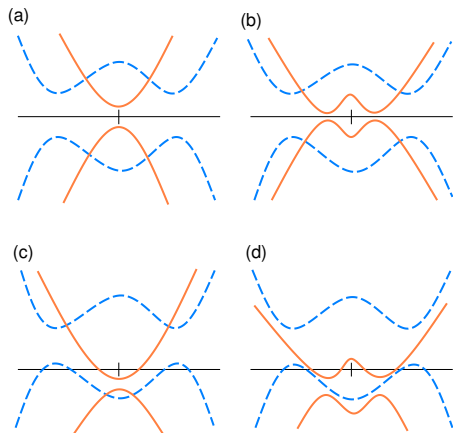
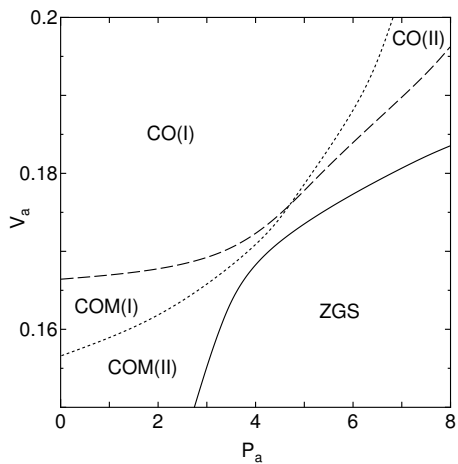
$$H_\sigma(\mathbf{k}) = \sum_{\alpha\beta} \tilde{\epsilon}_{\alpha\beta\sigma}(\mathbf{k}) a_{\mathbf{k}\alpha\sigma}^\dagger a_{\mathbf{k}\beta\sigma}$$

$$\tilde{\epsilon}_{\alpha\beta\sigma}(\mathbf{k}) = I_{\alpha\sigma} \delta_{\alpha\beta} + \epsilon_{\alpha\beta}(\mathbf{k})$$

$$I_{\alpha\sigma} = U_\alpha \langle n_{\alpha-\sigma} \rangle + \sum_{\beta'\sigma'} V_{\alpha\beta'} \langle n_{\beta'\sigma'} \rangle$$

$$\epsilon_{\alpha\beta}(\mathbf{k}) = \sum_{\delta} t_{\alpha\beta} e^{i\mathbf{k}\cdot\delta}$$

Phase diagram

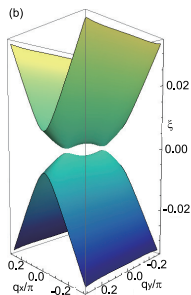
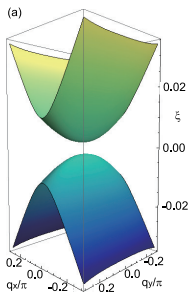


Effective 2-band Hamiltonian

$$H^{LK}(\mathbf{k}_M + \mathbf{q}) = f_0(\mathbf{q})\sigma_0 + f_1(\mathbf{q})\sigma_1 + f_2(\mathbf{q})\sigma_2 + f_3(\mathbf{q})\sigma_3$$

$$\xi_{\pm}(\mathbf{q}) = f_0(\mathbf{q}) \pm \sqrt{f_1(\mathbf{q})^2 + f_2(\mathbf{q})^2 + f_3(\mathbf{q})^2}$$

$$\Delta(\mathbf{q}) = \frac{\xi_+(\mathbf{q}) - \xi_-(\mathbf{q})}{2} \simeq \Delta + \frac{1}{2}(q_x \ q_y) S_M \begin{pmatrix} q_x \\ q_y \end{pmatrix} + \dots$$

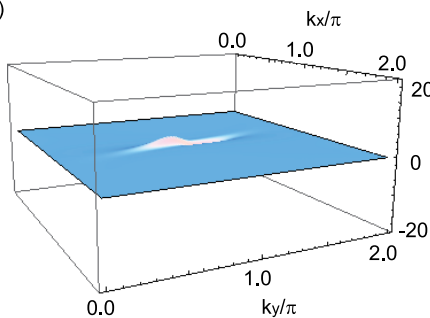


Berry curvature associated with the topological transition

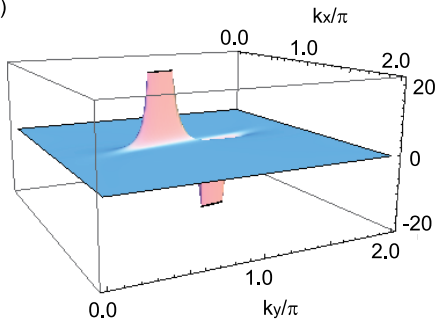
$$B_{n\sigma}(\mathbf{k}) = -i \sum_{m \neq n} \frac{v_{nm\sigma}^x(\mathbf{k})v_{mn\sigma}^y(\mathbf{k}) + c.c.}{(\xi_{n\sigma}(\mathbf{k}) - \xi_{m\sigma}(\mathbf{k}))^2}$$

$$v_{nm\sigma}^{x,y}(\mathbf{k}) = \langle n\sigma(\mathbf{k}) | \partial_{k_{x,y}} H_{\sigma}(\mathbf{k}) | m\sigma(\mathbf{k}) \rangle = \sum_{\alpha\beta} d_{n\alpha\sigma}(\mathbf{k})^* d_{m\beta\sigma}(\mathbf{k}) \partial_{k_{x,y}} \epsilon_{\alpha\beta}(\mathbf{k})$$

(a)



(b)



Berry phase and Chern number

$$Ch_{1\uparrow} = \frac{1}{2\pi} \int_{BZ} dSB_{1\uparrow}(\mathbf{k}) = 0$$

$$\Gamma(\mathbf{k}_{\pm}) = \frac{1}{2\pi} \int_{S(\mathbf{k}_{\pm})} dSB_{1\uparrow}(\mathbf{k}) \neq 0$$

