

Majorana fermions in superconducting nanowires without spin-orbit coupling

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We show that confined Majorana fermions can exist in nanowires with proximity induced *s*-wave superconducting pairing if the direction of an external magnetic field rotates along the wire. The system is equivalent to nanowires with Rashba-type spin-orbit coupling, with strength proportional to the derivative of the field angle. For realistic parameters, we demonstrate that a set of permanent magnets can bring a nearby nanowire into the topologically non-trivial phase with localized Majorana modes at its ends. Without the requirement of spin-orbit coupling this opens up for a new route for demonstration and design of Majorana fermion systems.

Journal Club

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Majorana bound states in nanowires

Ingredients :

- ▶ Semiconductor with strong Rashba SOI (e.g. InAs)
 - ▶ Zeeman field perpendicular to SOI
 - ▶ Proximity induced superconductivity
- } Helical conductor

$$\mathcal{H} = [p^2/2m - \mu(y)]\tau_z + up\sigma_z\tau_z + B(y)\sigma_x + \Delta(y)\tau_x$$

1D topological superconductor for

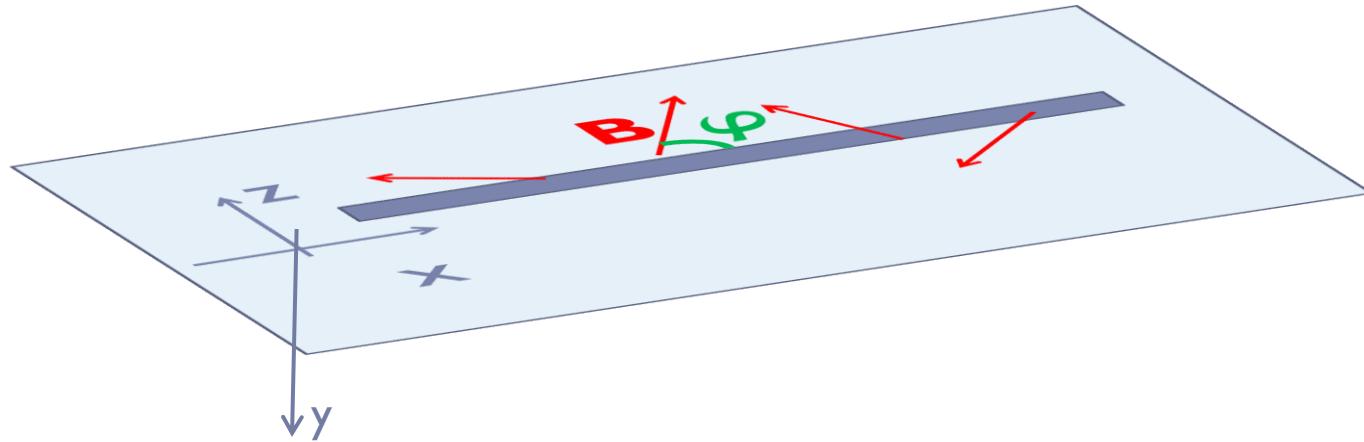
$$B^2 > \Delta^2 + \mu^2$$



Majorana bound states at the interfaces

- ▶ Oreg, Refael, von Oppen, PRL 105, 177002 (2010)

Engineering an effective Rashba SOI



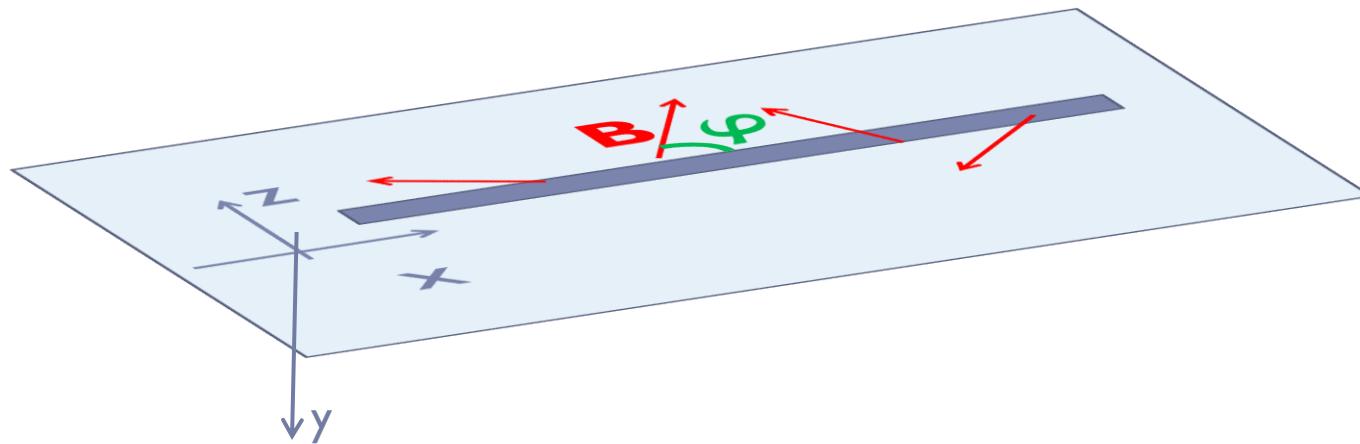
Inhomogeneous field

$$H = \frac{p^2}{2m} + \frac{1}{2}g\mu_B B(x) \cdot \sigma$$

Rotate locally the spin basis $U = e^{i\sigma_y \varphi/2}$



Engineering an effective Rashba SOI



Inhomogeneous field

$$H = \frac{p^2}{2m} + \frac{1}{2}g\mu_B B(x) \cdot \sigma$$

Rotate locally the spin basis $U = e^{i\sigma_y \varphi/2}$

$$\tilde{H} = \frac{p^2}{2m} + \frac{\hbar^2 \varphi'^2}{2m} + \frac{\hbar^2 \varphi''}{4m} \sigma_y + \frac{1}{2}g\mu_B B \sigma_z + \frac{\hbar \varphi'}{2m} p \sigma_y$$

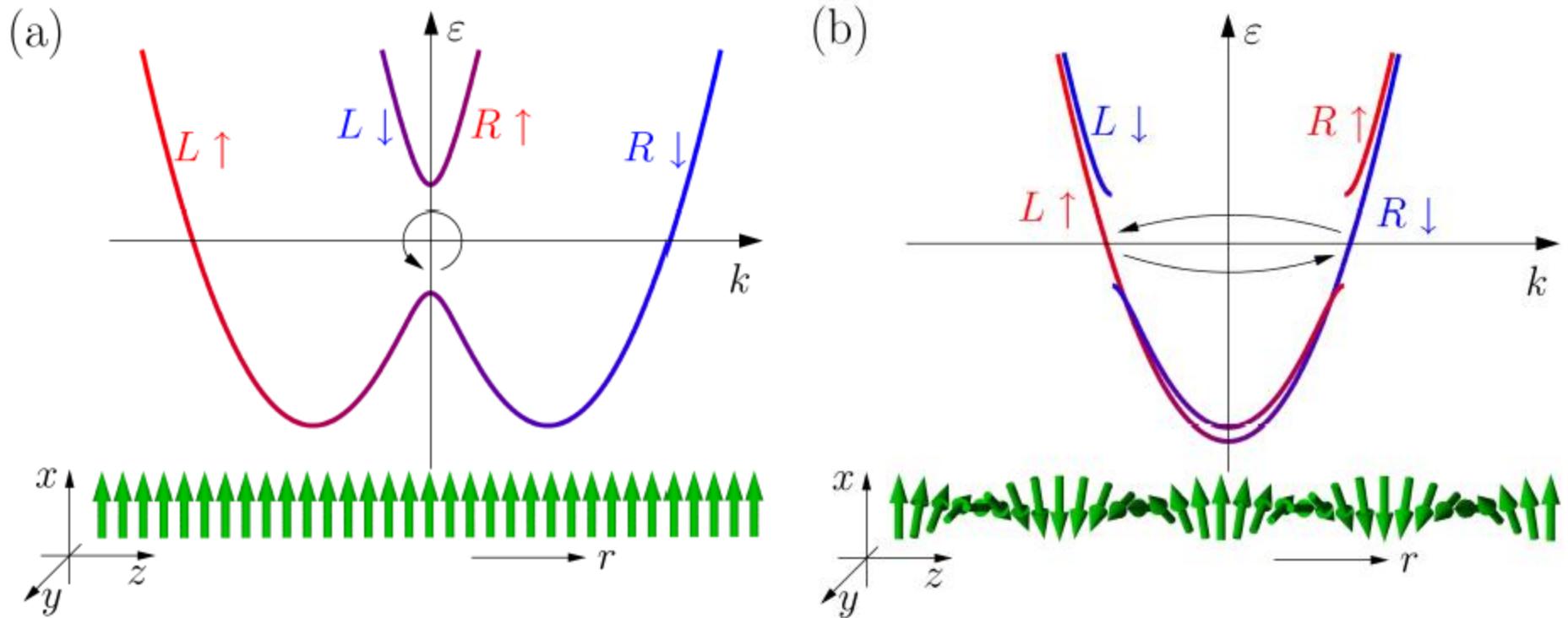
Rashba SOI!

Renormalizes μ

Effective Zeeman field



Rashba SOI == helical magnetic field



$$\tilde{H} = \frac{p^2}{2m} + \frac{\hbar^2\varphi'^2}{2m} + \frac{\hbar^2\varphi''}{4m}\sigma_y + \frac{1}{2}g\mu_B B\sigma_z + \frac{\hbar\varphi'}{2m}p\sigma_y$$

Side effect : effective Zeeman field is tilted toward SOI field

Optimal case

$$B = B_0(\cos(x/R), 0, \sin(x/R)) \quad \varphi' = 1/R \quad \varphi'' = 0$$

For low mass material $m = 0.014 m_e$ (InSb) and $R = 100$ nm, strength of SOI is

$$\alpha_{\text{eff}} = \frac{\hbar}{2mR} \approx 3 \times 10^4 \text{ m/s}$$

(similar to InAs)

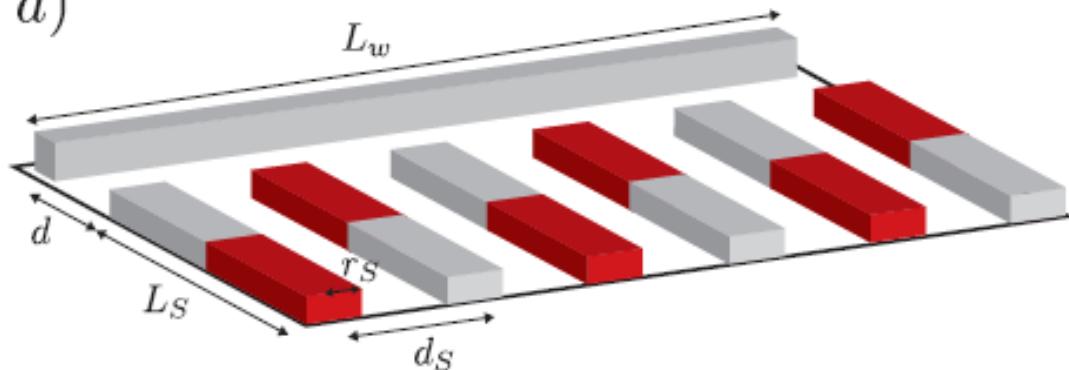
Topological phase for

$$g^*\mu_B|B_c| > \sqrt{|\Delta|^2 + (\mu - \hbar^2/8mR^2)^2}$$



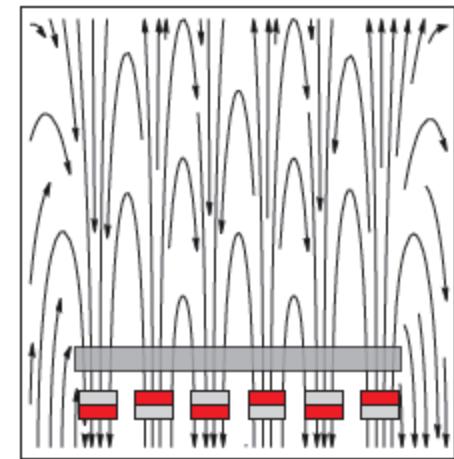
Anti-parallel permanent micro magnets

a)

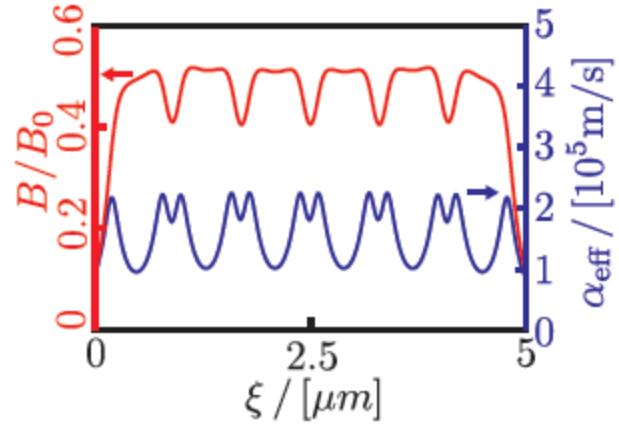


$$\begin{aligned}L_w &= 5 \mu\text{m} \\d &= 100 \text{ nm} \\L_s &= 600 \text{ nm} \\r_s &= 330 \text{ nm} \\d_s &= 860 \text{ nm}\end{aligned}$$

a)



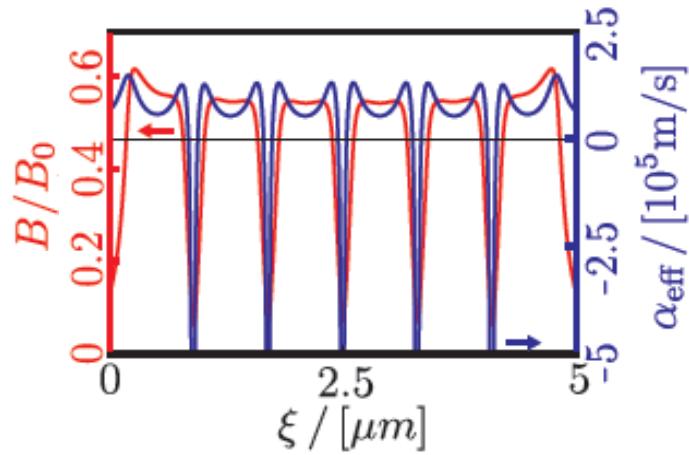
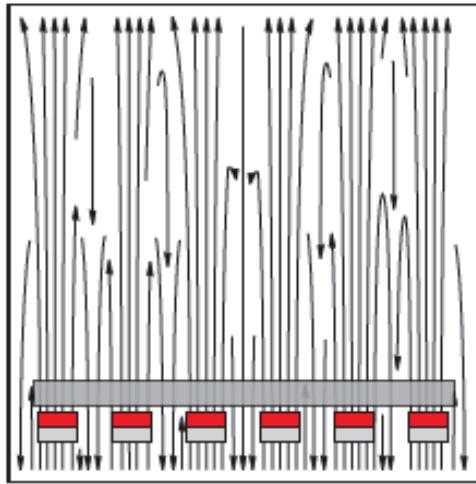
b)



Realization : magnets with different hysteresis loops



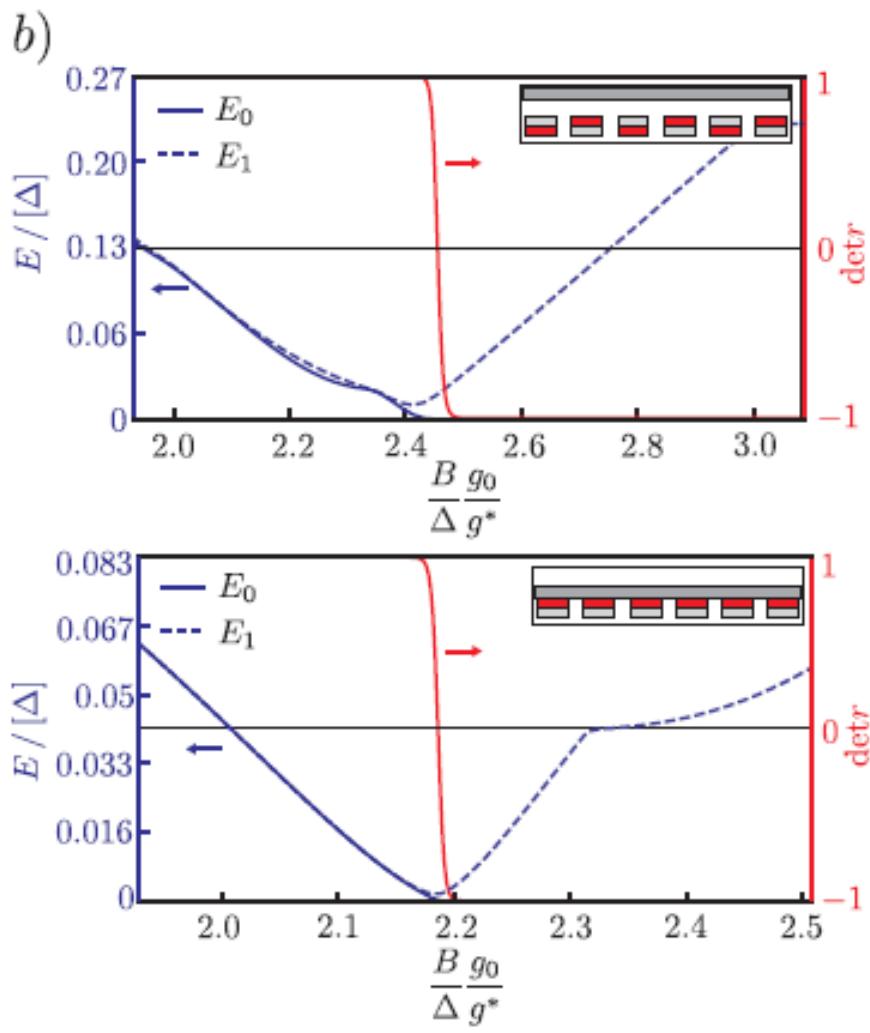
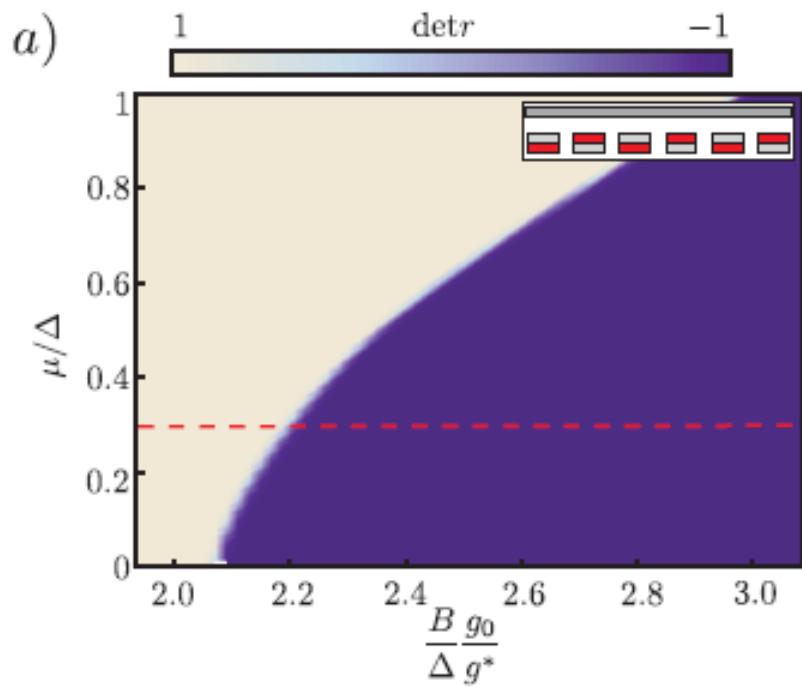
Parallel configuration



SOI and Zeeman go through zero;
Does the topological phase survive ?

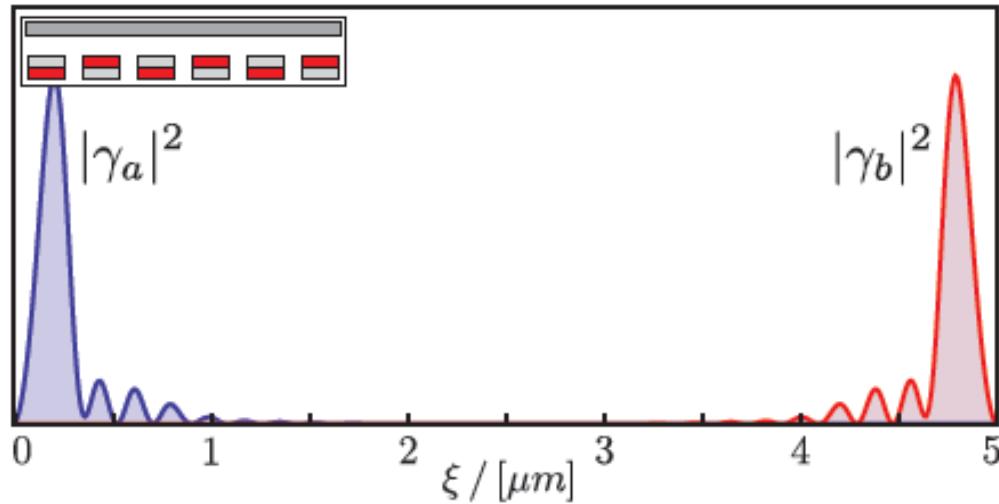


Phase diagram

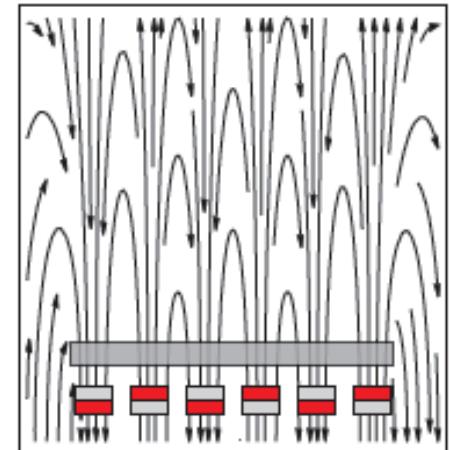
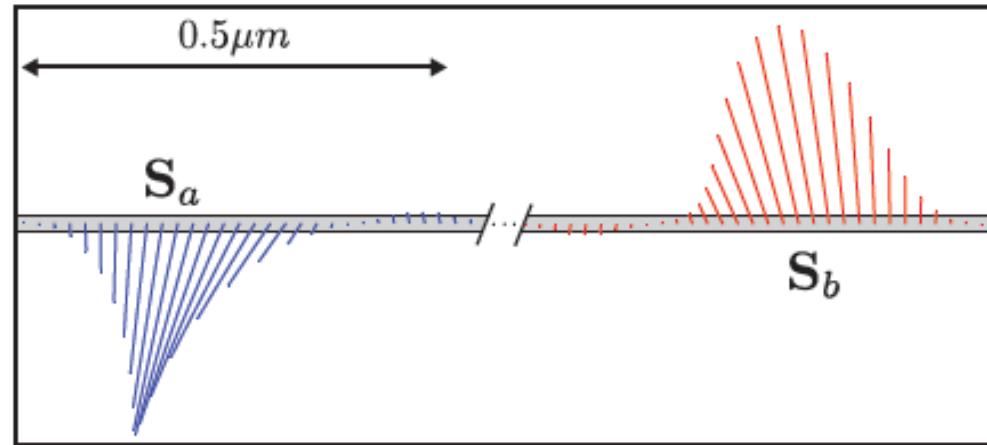


Structure of the Majorana bound states

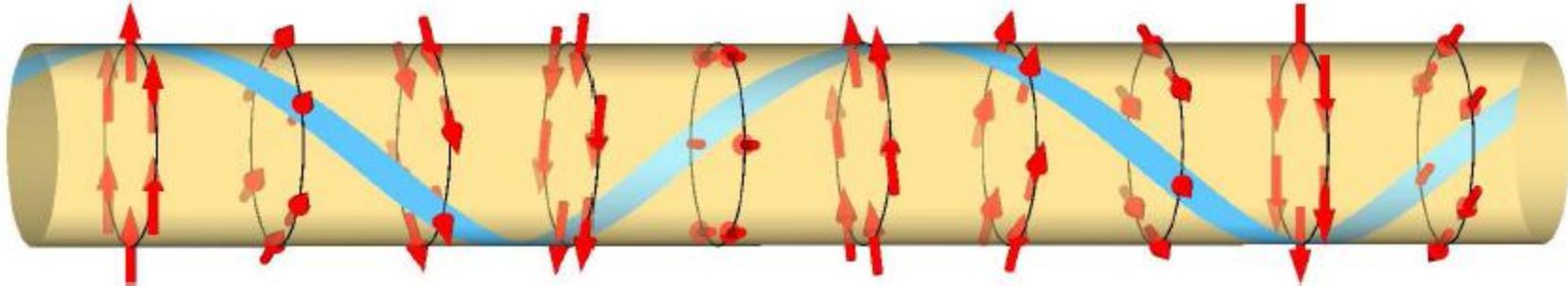
a)



b)



Helical Overhauser field



Braunecker, Simon, Loss, PRB **80**, 165119 (2009)

Already suggested as a candidate for Majorana fermions :

Gangadharaiah, Braunecker, Simon, Loss, PRL **107**, 036801 (2011)



Conclusion

- ▶ Helical magnetic field can replace SOI
- Focus on materials optimizing other properties (e.g. g-factor)

