

**Unbounded growth of entanglement  
in models of many-body localization**

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**arXiv:1202.5532**

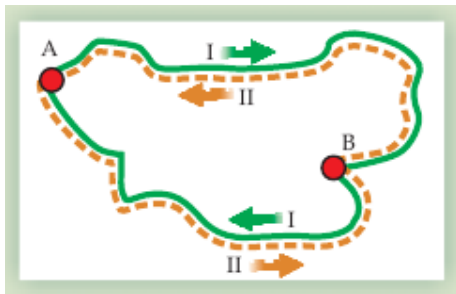
Vladimir M. Stojanović @ Journal Club

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## Anderson localization: the basics

“Absence of diffusion in certain random lattices”,  
P. W. Anderson, Phys. Rev. **109**, 1492 (1958)

a wave-packet (or a particle) moving in a spatially-disordered,  
( time-independent ) potential exhibits localization!



After: Legendijk et al., Phys. Today (2009)

A. L. most easily demonstrated in optical- and matter-wave systems

## What is this paper about?

**Q:** Can a closed quantum system of many interacting particles be localized by disorder?

**Recent studies:** in the absence of other degrees of freedom (e.g., phonons), e-e interaction may give rise to a “many-body localization transition” even in 1d!

**Goal of this paper:** show that the “many-body localized” phase differs qualitatively from the conventional (non-interacting) localized phase even in 1d!

**Outcome:** entanglement entropy and particle-number fluctuation show slow logarithmic evolution in time!  
entanglement entropy does not saturate in the thermodynamic limit, even for very weak interactions!

## Random-field ( $s = 1/2$ ) $XX$ Hamiltonian in $1d$

$$H_0 = J_{\perp} \sum_{i=1}^{N-1} \left( S_i^x S_{i+1}^x + S_i^y S_{i+1}^y \right) + \sum_{i=1}^N h_i S_i^z$$

$h_i$  – uniform random numbers  $\in [-\eta, \eta]$

**Reminder:** the Jordan-Wigner (JW) transformation

$$S_1^- = c_1 \quad , \quad S_i^- = \exp[i\pi \sum_{l=1}^{i-1} c_l^\dagger c_l] c_i \quad (i \geq 2)$$

maps  $H_0$  onto a Hamiltonian for free fermions with  
n.n. hopping and random on-site potential ( $S_i^z = c_i^\dagger c_i - 1/2$ )

study time-evolution of an initially unentangled pure state

the von Neumann entropy:  $S = -\text{Tr}_A(\hat{\rho}_A \ln \hat{\rho}_A) = -\text{Tr}_B(\hat{\rho}_B \ln \hat{\rho}_B)$

bipartition  $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$  by dividing the system at the center bond!

for asymptotic behavior ( $t \rightarrow \infty$ ) use exact diagonalization:

$$H = \sum_j \lambda_j P_j \quad \Longrightarrow \quad U(t) = \sum_j e^{-i\lambda_j t} P_j$$

results obtained by averaging over  $> 10^4$  field configurations  $\{h_i\}$ ,  
starting from a random state  $|\Psi_{t=0}\rangle = |m_1\rangle \otimes |m_2\rangle \dots \otimes |m_L\rangle$   
 $m_j \in \{\uparrow, \downarrow\}$

## Adding the $z$ (Ising) coupling (fermion interactions)

$$H = \sum_{i=1}^{N-1} \left[ J_{\perp} \left( S_i^x S_{i+1}^x + S_i^y S_{i+1}^y \right) + J_z S_i^z S_{i+1}^z \right] + \sum_{i=1}^N h_i S_i^z$$

**Q:** Is there any qualitative change in the behaviour of physical quantities when a small interaction ( $J_z$ ) is added?

**Already known (almost...problem to approach the TD limit!)**

V. Oganesyan and D. Huse, PRB **75**, 155111 (2007)

for  $J_{\perp} = J_z$  (Heisenberg case) the system undergoes a dynamical transition as a function of  $\eta/J_z$  (in all eigenstates):

for small enough  $\eta/J_z$  (strong-enough interactions)

localization is destroyed (i.e., spin conductivity is nonzero)!

## Entanglement growth after a quench

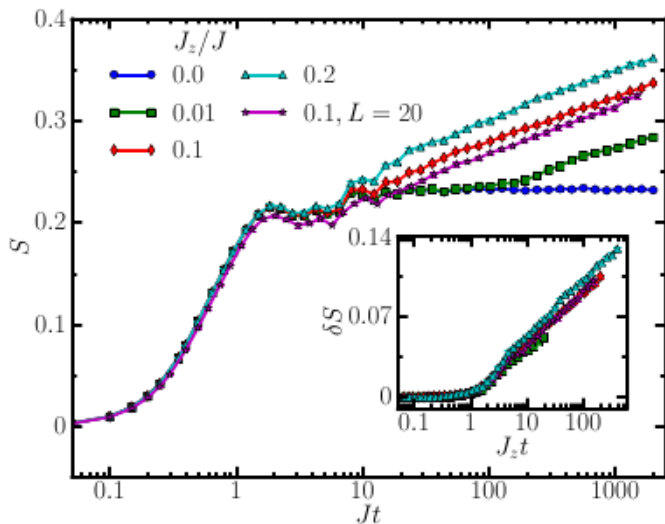
for  $J_z = 0$  (no interactions) the “half-chain” entanglement entropy saturates; saturation sets in at times  $\sim J_{\perp}^{-1}$

**expectation:** ...weak interaction leads to a small delay in saturation and a small increase in final entanglement...

instead, entanglement growth shows a qualitative change of behavior already for infinitesimal  $J_z$ !

It grows logarithmically even after times  $\gg J_{\perp}^{-1}$ !

# Entanglement growth after a quench

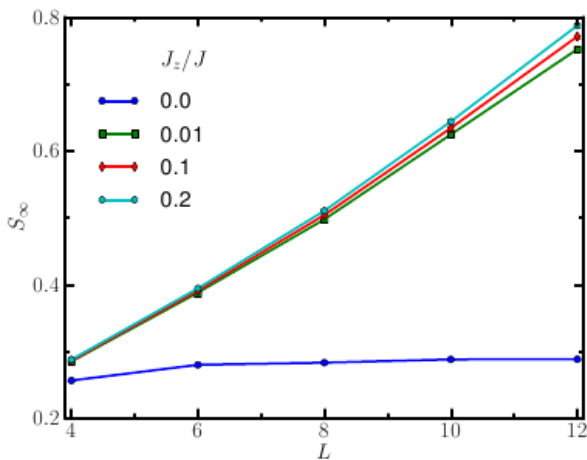


data for  $\eta = 5$  and  $L = 10$



## Saturation value of entropy as a function of the system size

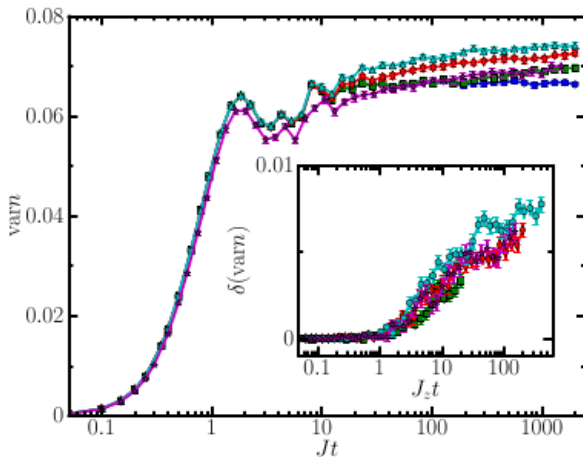
saturation values for finite  $L$  are essentially independent of  $J_z$ !



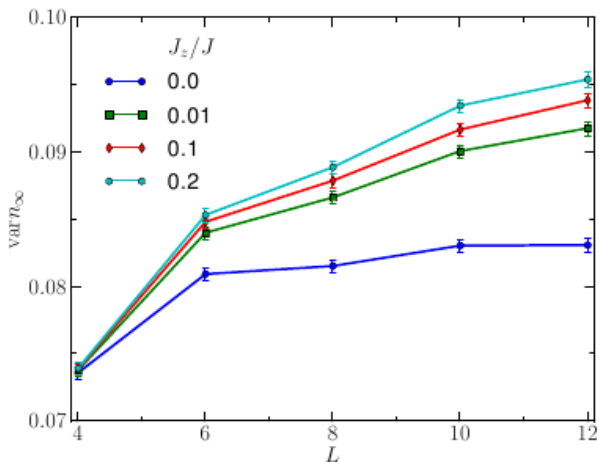
the entanglement entropy does not saturate in the TD limit!

# Half-chain particle number fluctuations

i.e., the variance of the total spin on half the chain



# Saturation values of the half-chain particle-number variance



unlike for entanglement, saturation values depend on the interaction strength!