

# To close or not to close: the fate of the superconducting gap across the topological quantum phase transition in Majorana-carrying semiconductor nanowires

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arXiv:1206.0013

## Class D spectral peak in Majorana quantum wires

Dmitry Bagrets and Alexander Altland  
arXiv:1206.0434

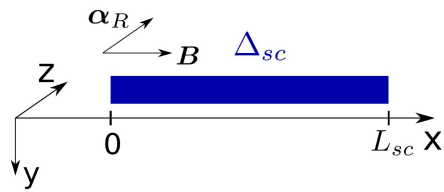
## Enhanced zero-bias Majorana peak in disordered multi-subband quantum wires

Falko Pientka, Graham Kells, Alessandro Romito, Piet W. Brouwer, and Felix von Oppen  
arXiv:1206.0723

## Zero-bias peaks in spin-orbit coupled superconducting wires with and without Majorana states

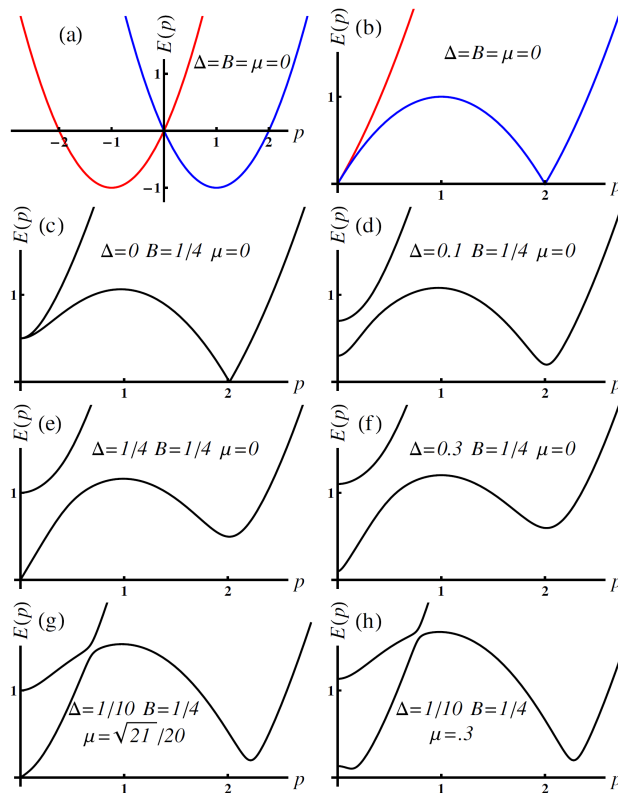
Jie Liu, Andrew C. Potter, K.T. Law, and Patrick A. Lee  
arXiv:1206.1276

# Majorana Fermions in Semiconducting Nanowires



$$H = \int \Psi^\dagger(y) \mathcal{H} \Psi(y) dy; \quad \Psi^\dagger = (\psi_\uparrow^\dagger, \psi_\downarrow^\dagger, \psi_\downarrow, -\psi_\uparrow)$$

$$\mathcal{H} = [p^2/2m - \mu(y)]\tau_z + up\sigma_z\tau_z + B(y)\sigma_x + \Delta(y)\tau_x.$$



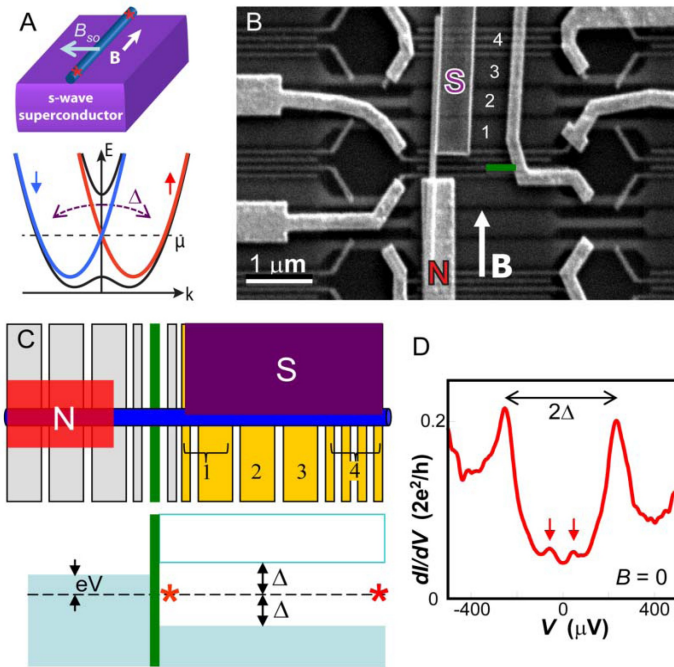
$$E_0 = E(p = 0) = |B - \sqrt{\Delta^2 + \mu^2}|$$

topological gap

$$B^2 > \Delta^2 + \mu^2$$

topological criterion

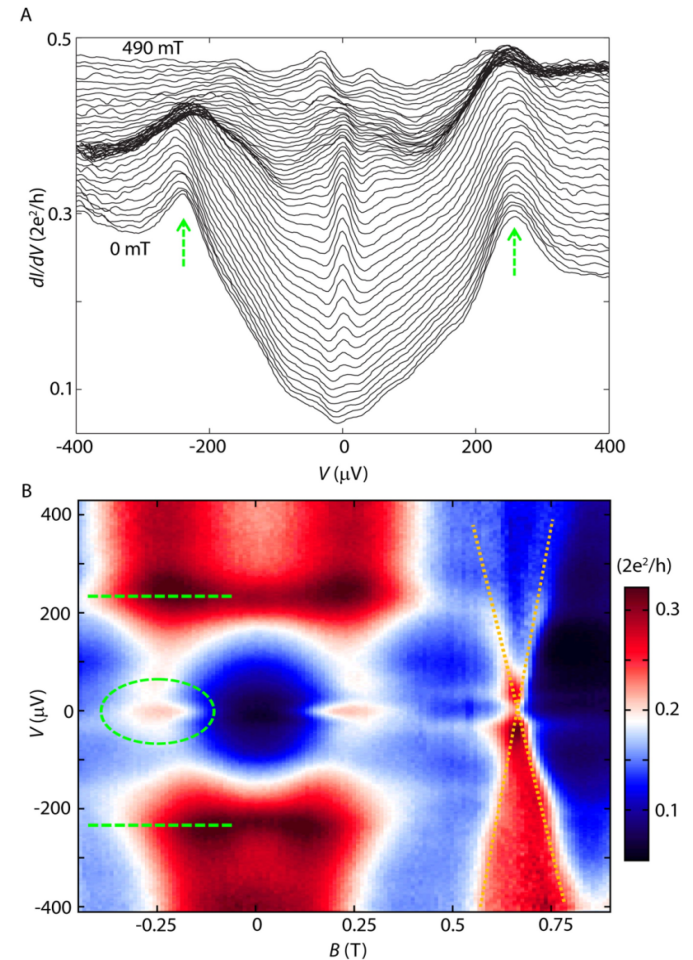
# Delft experiment



$$l_{\text{so}} \approx 200 \text{ nm}$$

$$\alpha \approx 0.2 \text{ eV} \cdot \text{\AA}$$

$$\Delta \approx 250 \text{ } \mu\text{eV}$$



no closing of the topological gap!!!

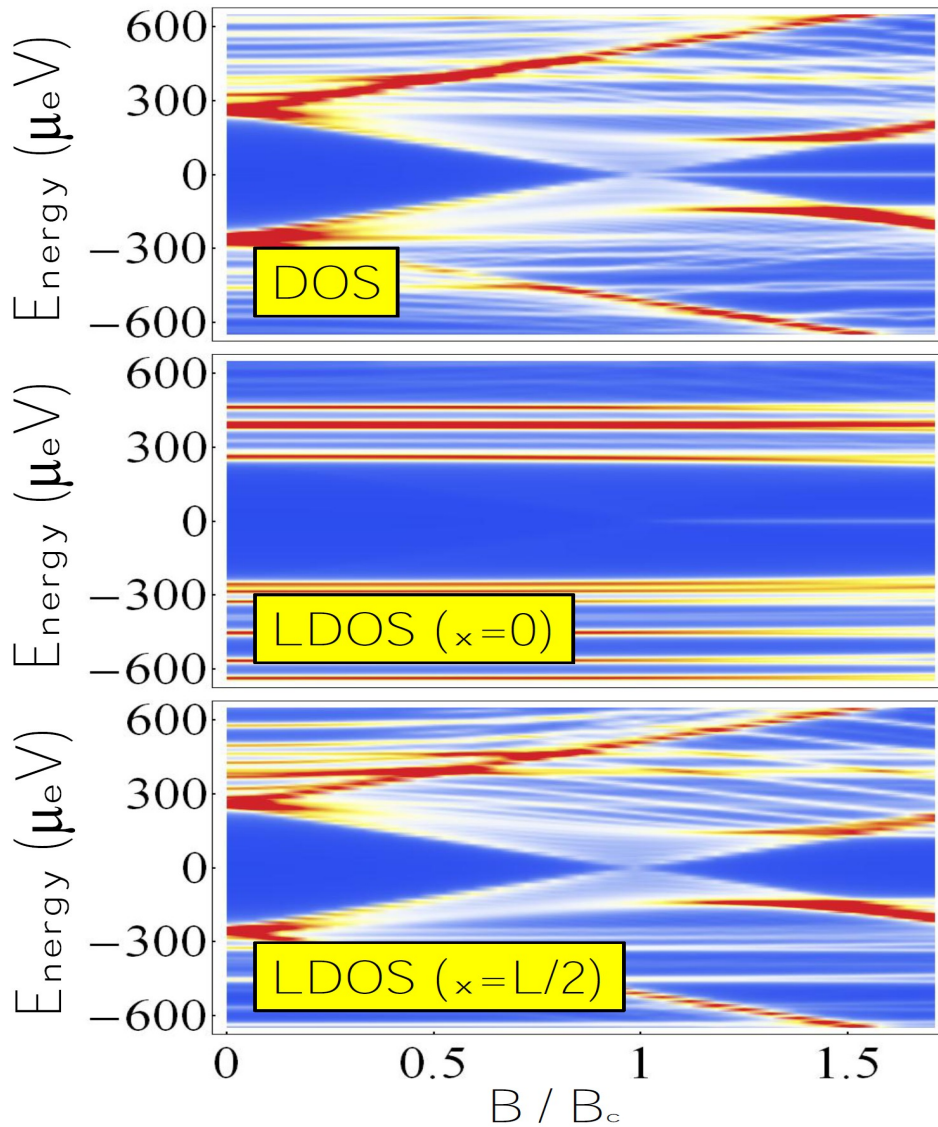
A rectangular semiconducting nanowire:  $L_x \gg L_y \sim L_z$   $(u_\uparrow, u_\downarrow, v_\uparrow, v_\downarrow)$

$$H_{nm}(k) = [\epsilon_{nm}(k) - \mu\delta_{nm}]\tau_z + \Gamma\delta_{nm}\sigma_x\tau_z + \alpha k\delta_{nm}\sigma_y\tau_z - i\alpha_y q_{nm}\sigma_x + \Delta_{nm}\sigma_y\tau_y,$$

chemical potential Zeeman term  $\Gamma = g^*\mu_B B/2$   
spin orbit interaction superconductivity  
 $\alpha = 0.2 \text{ eV\AA}$

$\sigma_i$  spin  
 $\tau_i$  particle-hole

modes:  
 $n = (n_y, n_z) \quad m = (m_y, m_z)$   
 $\sin(n_y\pi y/L_y) \sin(n_z\pi z/L_z)$



$$B = B_c \approx 0.2 \text{ T}$$

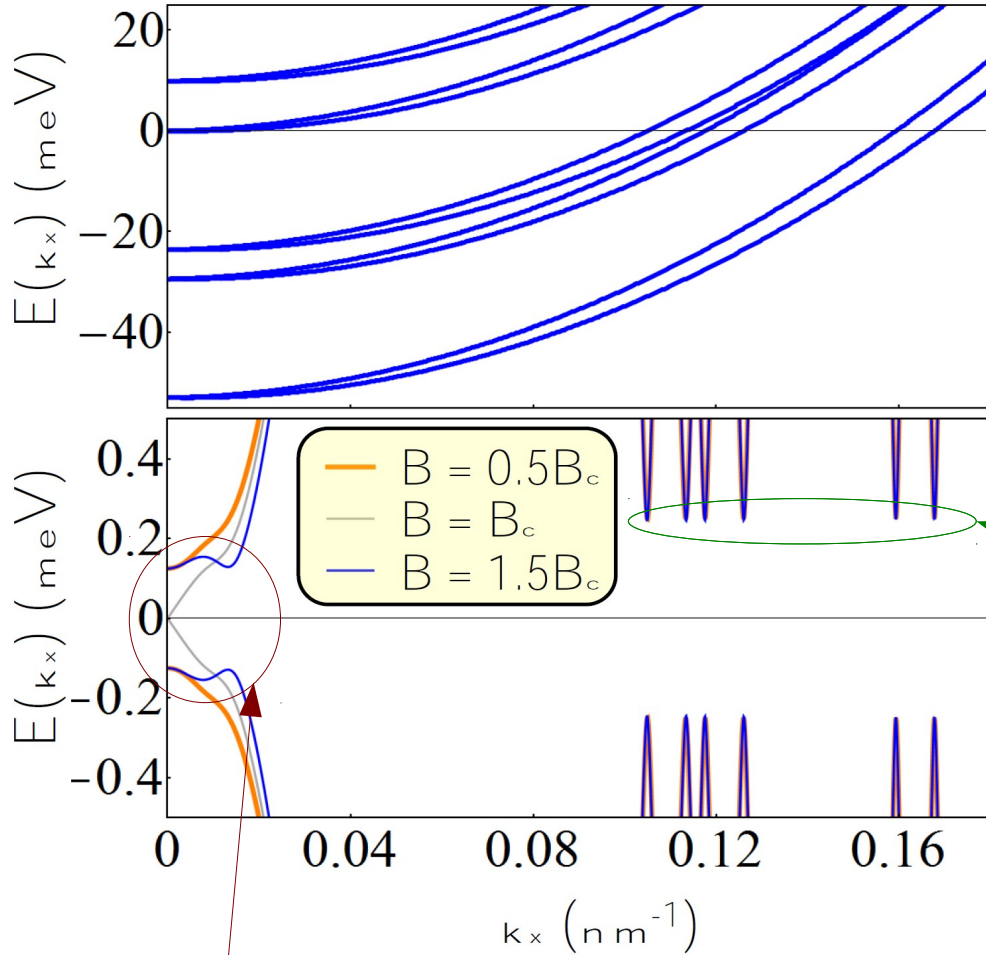
$$\mu = \Delta/2$$

$$\Delta \sim 250 \mu\text{eV}$$

Assumption: limit of point contact tunneling

*“The main difference between our calculated LDOS and the measured differential tunneling charge conductance current is some matrix element effects which are neither precisely known nor of any qualitative significance.”*

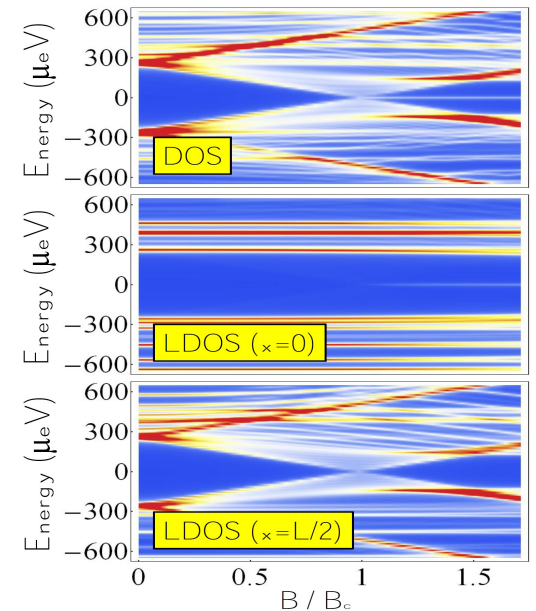
# LDOS vs. DOS

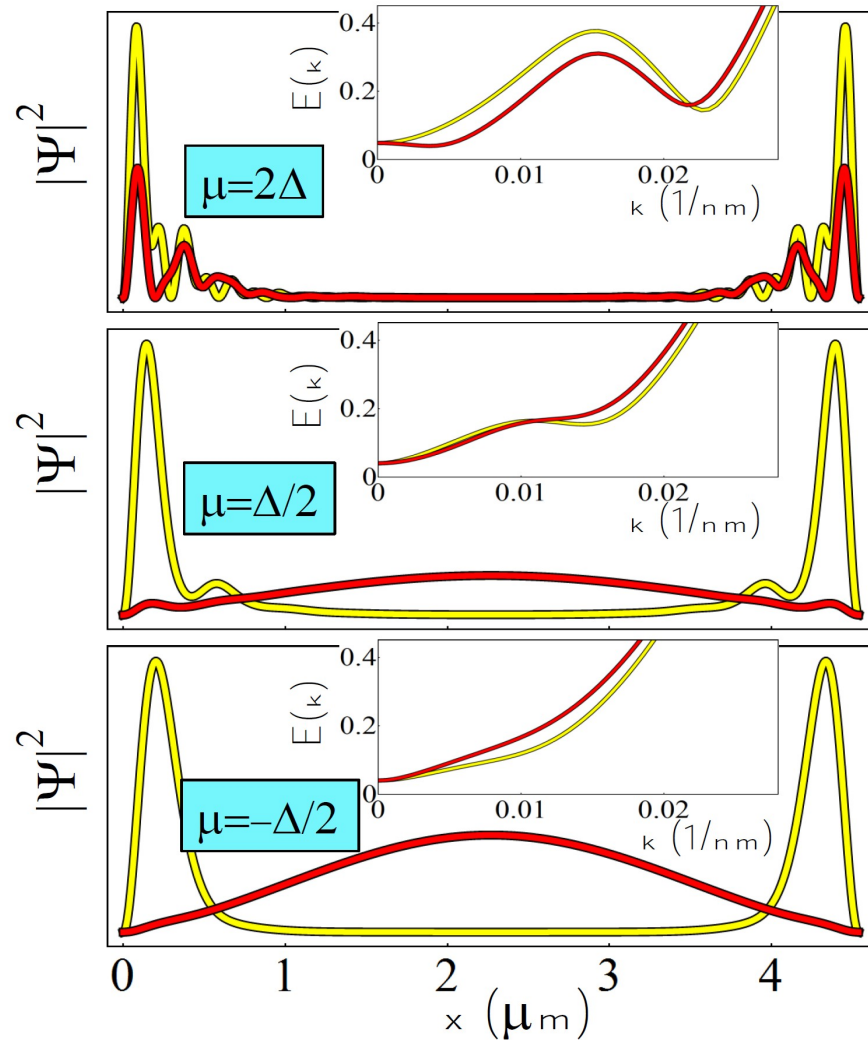


four occupied bands  $\mu = \Delta/2$

low-energy bands  
weakly depend on B

top band - "Majorana band"  
strongly depend on B  
topological gap vanishes and reopens





$$\underline{B = 0.9B_c}$$

non-topological

$$\underline{B = 1.1B_c}$$

topological

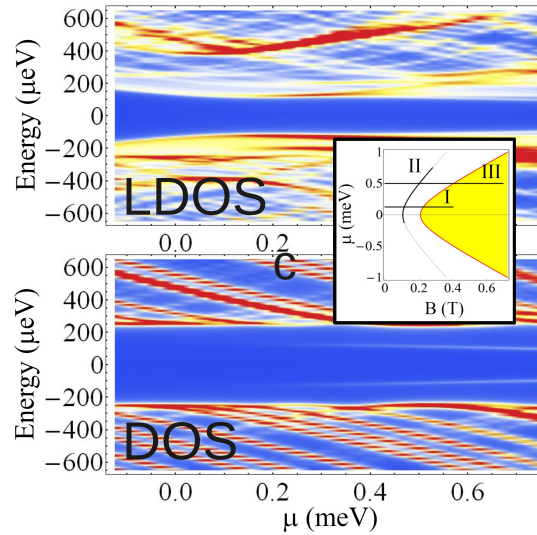
*“In the non-topological SC phase, we find that the lowest energy states are delocalized when the chemical potential is below a certain value”*

$$\mu_c(\Gamma) \sim \mathcal{O}(\Delta)$$

# LDOS vs. DOS

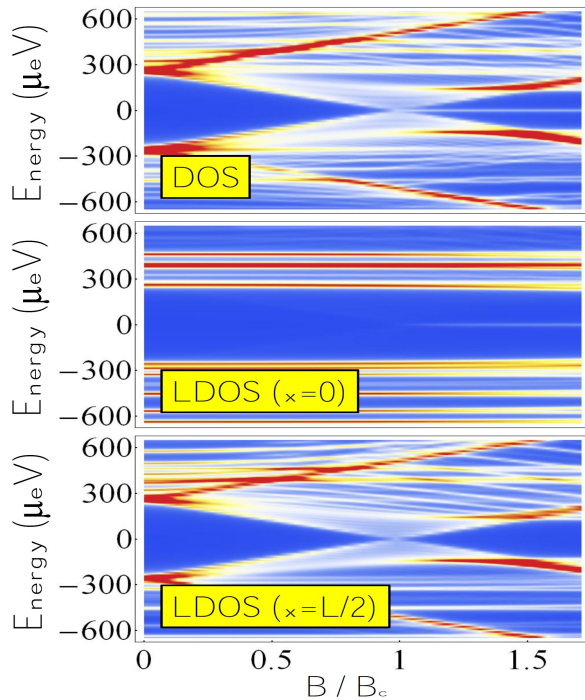
delocalized

$$\mu < \mu_c$$

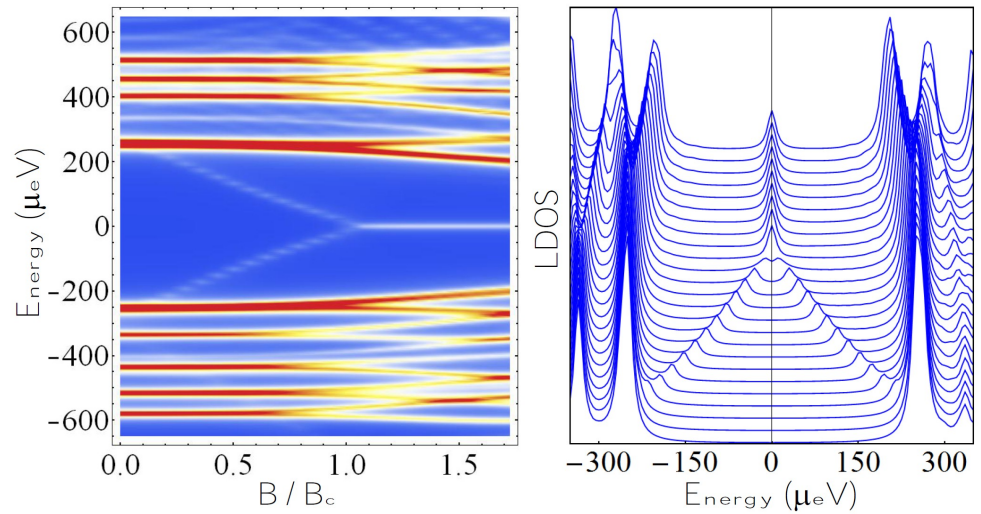


localized

$$\mu > \mu_c$$



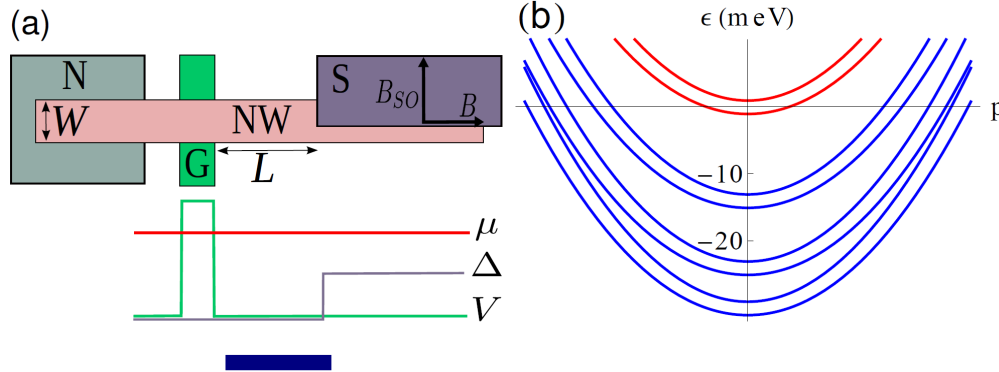
no closing of the topological gap is observed



closing of the topological gap is observed



# Enhanced zero-bias Majorana peak - disorder



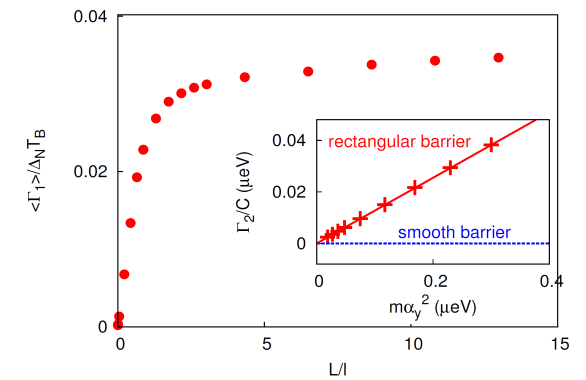
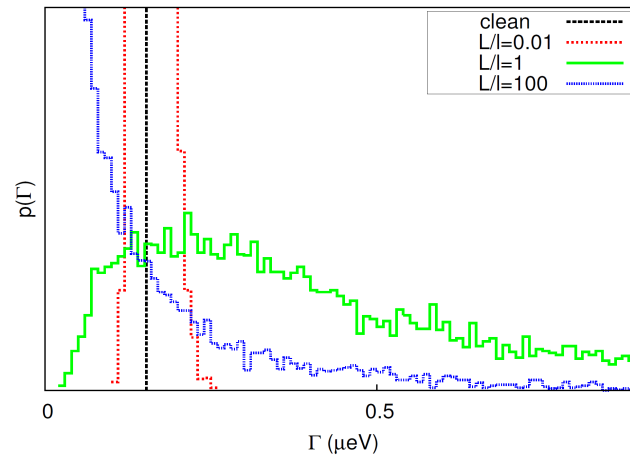
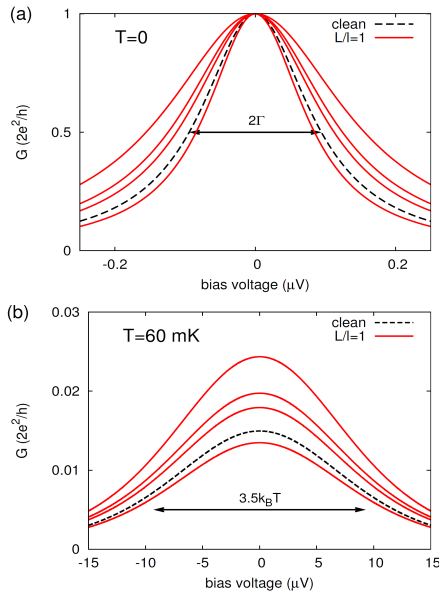
$$\mathcal{H} = \left( \frac{\mathbf{p}^2}{2m} + \alpha p_x \sigma_y - \alpha_y p_y \sigma_x + U(x) + V_{\text{dis}}(\mathbf{r}) - \mu \right) \tau_z - B \sigma_x + \Delta(x) \tau_x,$$

$$\langle V_{\text{dis}}(\mathbf{r}) V_{\text{dis}}(\mathbf{r}') \rangle = \frac{v_F^2}{l_{2d}} \delta(\mathbf{r} - \mathbf{r}')$$

disorder between the tunnel barrier and superconductor

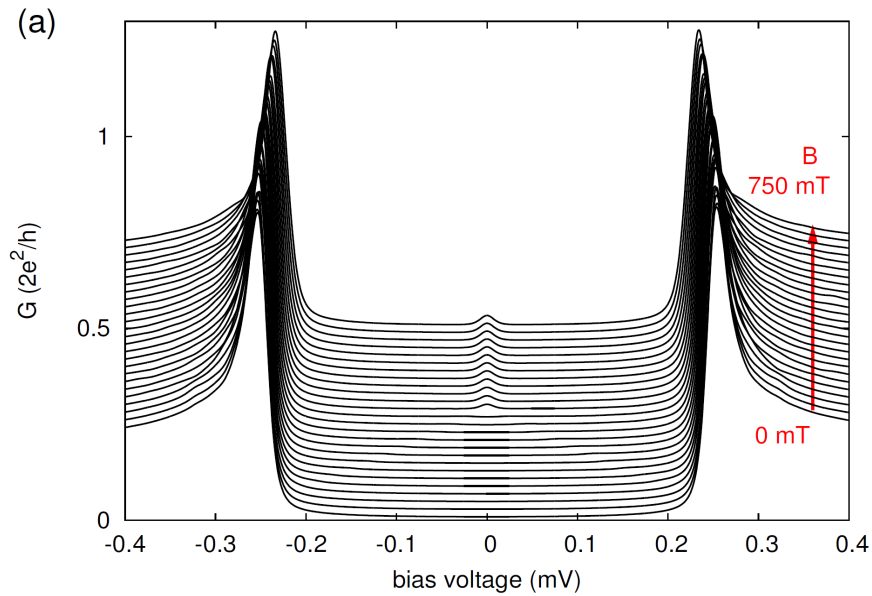
$$G(V) = \frac{e^2}{h} \text{tr} [1 + r_{\text{he}}(eV) r_{\text{he}}(eV)^\dagger - r_{\text{ee}}(eV) r_{\text{ee}}(eV)^\dagger]$$

## Zero-bias conductance peak

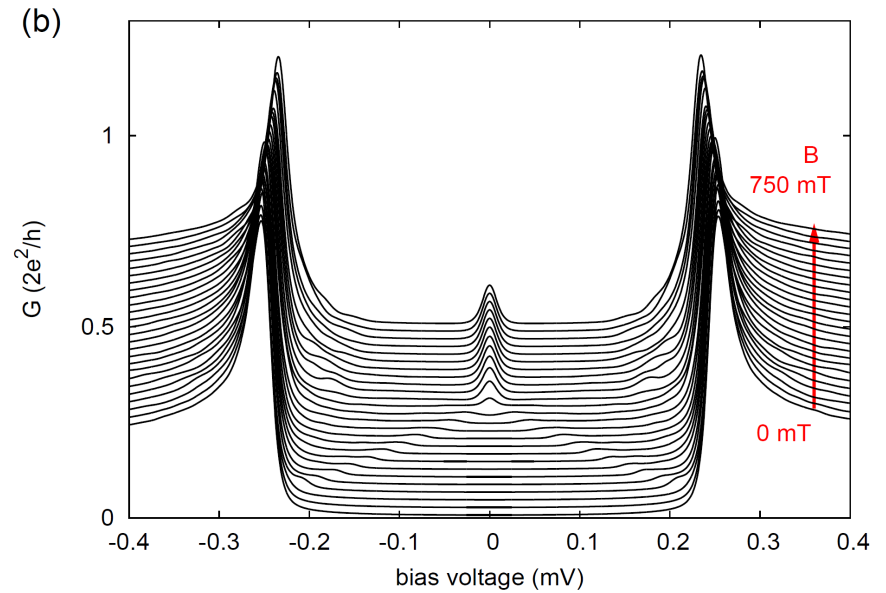


# Enhanced zero-bias Majorana peak - disorder

Piet Brouwer  
Felix von Oppen



clean wire

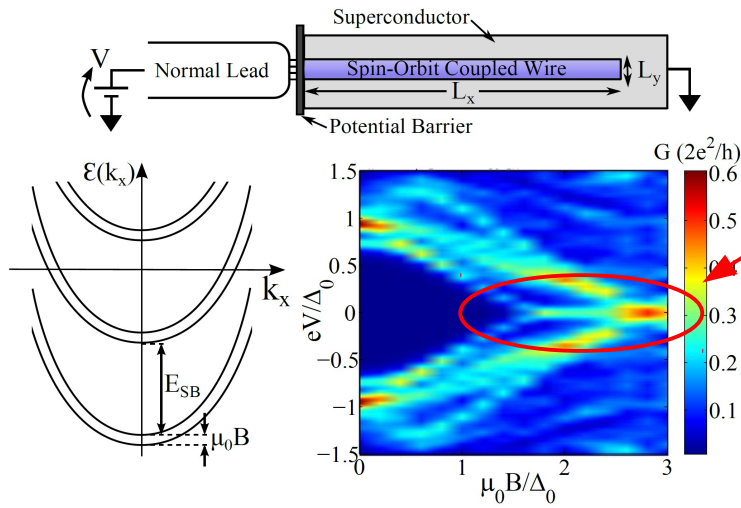


weak disorder

$$l = 10L$$

$$L = 10nm$$

# Zero-bias peaks with and without Majorana end-states



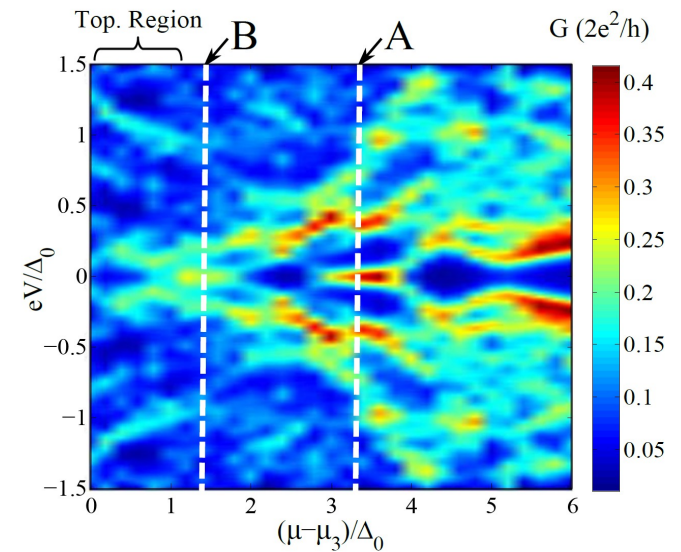
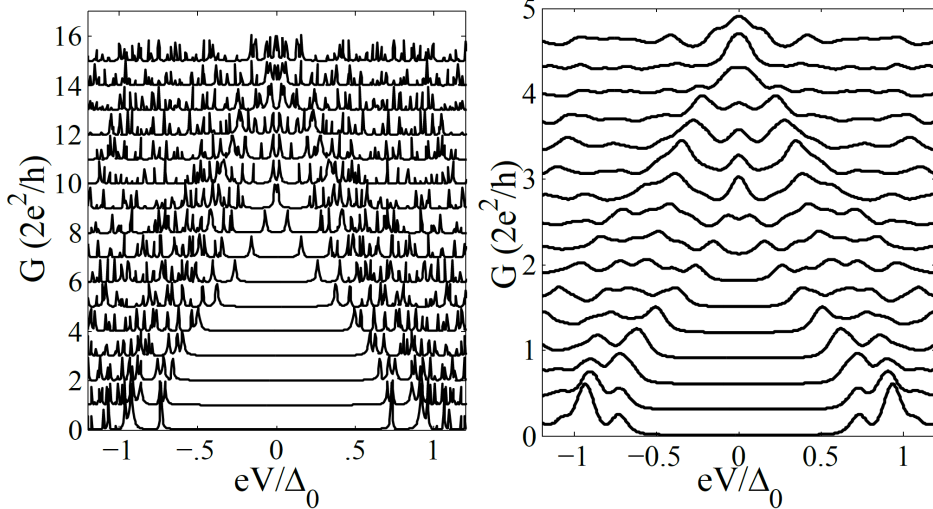
ZBP in the non-topological regime due to disorder

$$\ell \approx 3\mu\text{m}$$

$$\overline{V(\mathbf{r})V(\mathbf{r}')} = W^2\delta_{\mathbf{r},\mathbf{r}'}$$

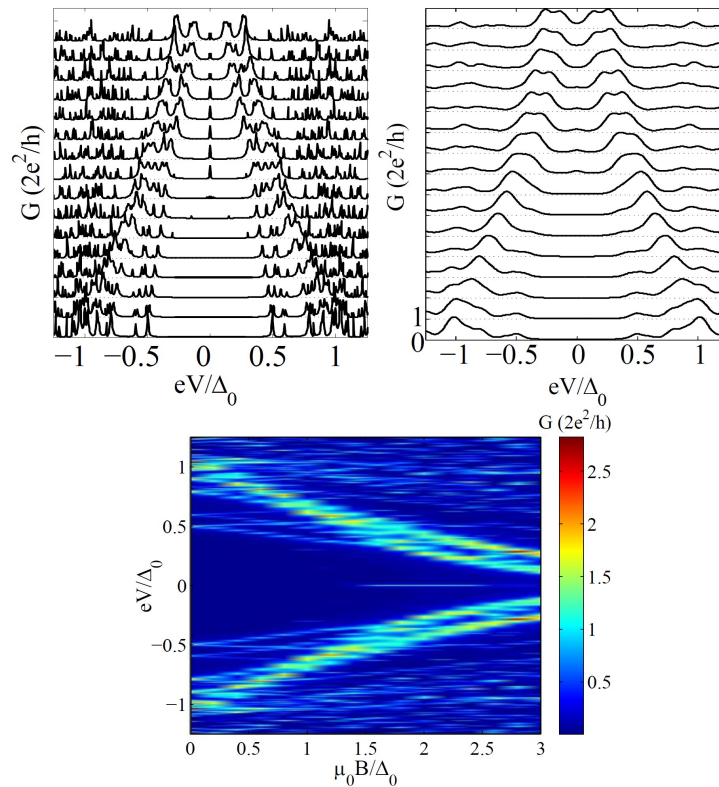
ZBP's exist in the non-topological state

ZBP's persist even when disorder is sufficiently strong



$$\ell \approx 10\mu\text{m}$$

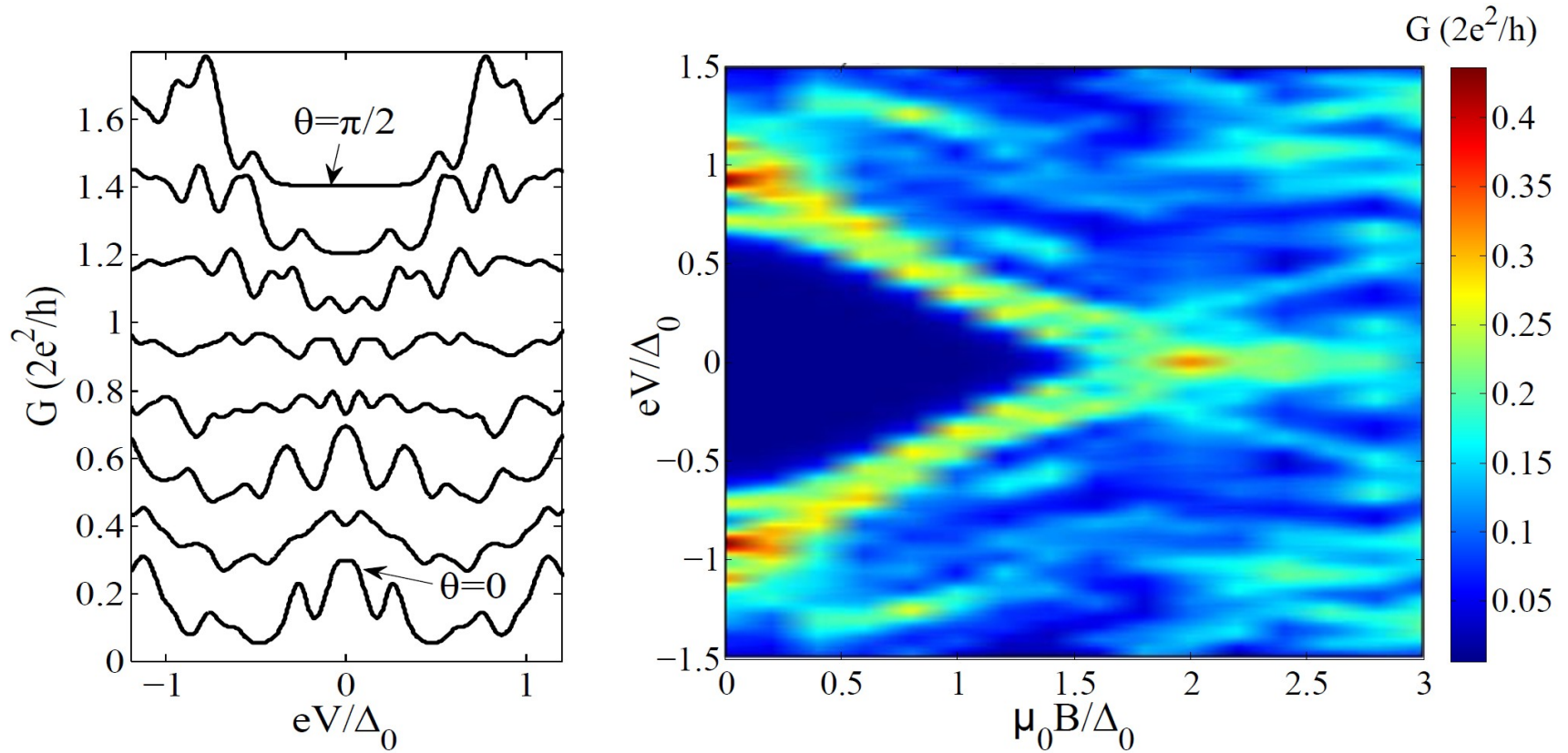
# Zero-bias peaks with and without Majorana end-states



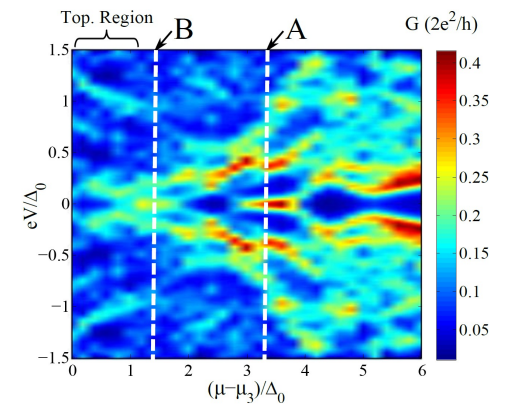
Long clean wire

*“We note that, even for a clean wire, the observation of this quantized peak requires much 2-3 times longer wires than those used in the experiment”*

# Zero-bias peaks with and without Majorana end-states



Angle dependence of non-topological ZBP



# Conclusion

*Jie Liu, Andrew C. Potter, K.T. Law, and Patrick A. Lee; arXiv:1206.1276*

- (i) Majorana end-states are destroyed and do not give rise to quantized ZBPs.
- (ii) ZBPs of a nontopological origin often appear due to disorder
- (iii) Non-topological ZBPs are typically stable
- (iv) Non-topological ZBPs appear and disappear under nearly identical conditions to those of true Majorana peaks.

*Dmitry Bagrets and Alexander Altland arXiv:1206.0434*

- (i) rigidly positioned at zero energy,
- (ii) strongly affected by temperature
- (iii) not affected by magnetic field,
- (iv) relies on parametric conditions similar to those required by the Majorana peak,
- (v) carries unit integrated spectral weight (a single quasi 'state'),
- (vi) independent of other midgap structures (Kondo resonances or Andreev bound states),
- (vii) shows sensitivity to the parity of channel number,  
unlike the Majorana particle,
- (viii) relies on the presence of a moderately weak amount of disorder

*“These results strike a note of caution for interpreting recent experimental evidence of Majorana states in tunneling data”*