# N. Sugimoto and N. Nagaosa: Spin-orbit echo Science **336**, 1413 (15 June 2012)

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#### The continuity equation

$$\partial_t \mathbf{p} = -\nabla \cdot \mathbf{j},$$
 giving

$$\partial_t Q = 0$$
,

is of no use for spin,

$$\partial_t \rho_s = -\nabla \cdot \mathbf{j}_s + \dots,$$

the spin-orbit coupling being among the dissipation sources.



The semiconductor 2DEG spin-orbit coupling,

$$H = \frac{1}{2m}\mathbf{p}^2 + V(\mathbf{x}) + \frac{\hbar}{2ml_{\alpha}}(\rho_x\sigma_y - \rho_y\sigma_x) + \frac{\hbar}{2ml_{\beta}}(-\rho_x\sigma_x + \rho_y\sigma_y),$$

can be written as

$$H = \frac{1}{2m}(\mathbf{p} - e\mathbf{A})^2 + V(\mathbf{x}),$$

with 
$$\mathbf{A} = \sum_i \mathbf{e}_i A_i^i \sigma_j$$
, and  $[A_\mu, A_\nu] = F_{\mu\nu} \neq 0$ .

## Alternative - the unitary basis transformation

(spin-orbit interaction "removal")

consider 
$$U = \exp(ie\mathbf{A} \cdot \mathbf{x}/\hbar)$$
, then

$$U^{\dagger}H(\mathbf{p},\mathbf{x})U=H(\mathbf{p}+\mathbf{e}\mathbf{A},\mathbf{x}),$$

so that you can remove the gauge from the kinetic energy

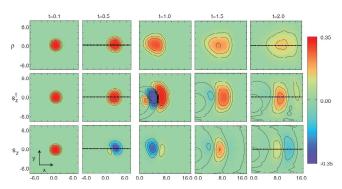
$$H_{\rm eff} = U^{\dagger} (H_0 + H_{\rm so}) U = H_0 + f(F_{\mu\nu}).$$

This is useless (due to *f*) except for:

- lacktriangleq quasi 0d system where  $F_{\mu\nu}$  is "small"  $(x,y\ll l_{so})$
- truly 1d system where  $F_{\mu\nu}$  is zero (single vector  $\mathbf{e}_i$ )
- 2d "compensated" system where  $F_{\mu\nu}$  is zero by chance  $(I_{\alpha} = I_{\beta})$

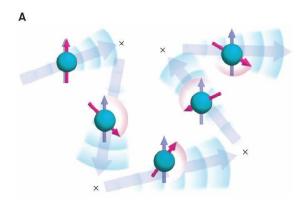
## Article main result 1: Twisted Spin

There exists coordinate frame where spin-orbit effects are trivial (=none) for general gauge-like spin-orbit interaction. Going back to a laboratoty frame, there is a conserved quantity.



#### Article main result 1 $\rightarrow$

In another words, Dyakonov-Perel relaxation is not really a relaxation (it is reversible, in principle).



# Special frame / Twisted Spin: $\mathbf{s}^t = (\hbar/2)\Psi^\dagger \Upsilon^\dagger \sigma \Upsilon \Psi$

$$\begin{split} & Y = \lim_{x' \to x} \frac{1}{2} \exp\{\frac{\mathrm{i}}{2} m(\alpha + \beta) \sigma_{-} x_{+}\} \times \left\{2 \sin^{2}\left[\frac{\mathrm{i}(\partial_{x} - \partial_{y})}{8 \sqrt{2} m(\alpha + \beta)}\right] + \right. \\ & \left. \exp\{-\frac{\mathrm{i}}{2} m(\alpha - \beta) \sigma_{+} x_{-}'\} \exp\{-\frac{\mathrm{i}}{2} \left[\frac{\partial_{x'}^{\leftarrow} \sigma_{z} \partial_{y} + \partial_{y'}^{\leftarrow} \sigma_{z} \partial_{x}}{2 m^{2} (\alpha^{2} - \beta^{2})}\right]\} \right. \\ & \left. + \exp\{-\frac{\mathrm{i}}{2} \left[\frac{\partial_{x'} \sigma_{z} \partial_{y} + \partial_{y'} \sigma_{z} \partial_{x}}{2 m^{2} (\alpha^{2} - \beta^{2})}\right]\} \exp\{-\frac{\mathrm{i}}{2} m(\alpha - \beta) \sigma_{+} x_{-}'\}\right\} \end{split}$$

#### properties:

- Y depends on electron velocity (reversible in principle, but hopeless in practice)
- Y is non-unitary
- lacksquare + twisted spin is really a spin:  $U_{SU(2)}^{\dagger} \mathbf{s} U_{SU(2)} = R_{SO(3)}[\mathbf{s}]$
- + persistent helix limit follows as a special case Y → exp(ieA·x/ħ)



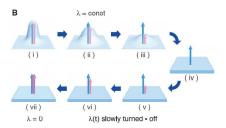
#### Article main result 2

There is a way to uncover the Twisted Spin: it is an adiabatic invariant.

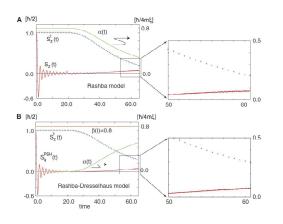
$$\partial_t \mathbf{s}^t = O\{(\partial_t I_{so}^{-1})^{n \ge 2}\}$$

turning off the spin-orbit strength adiabatically, the twisted spin is:

- 1. conserved,
- 2. reduced to the "standard" spin



# Spin-orbit echo numerical demonstration



- turning off SOC is not really possible; persistent helix will serve well
- adiabaticity requirement is in odds with inelastic spin dissipation channels (nuclei, EY, BAP, ee, ...) it is a question of material parameters now

# Derivation\* (supplementary information): preliminaries

- multiplication of two numbers  $\mu$  :  $a \otimes b \rightarrow a.b$  ( $\sim$ 3000BC)
- (Emmy) Noether theorem: To every 1-parametric (Sophus) Lie group symmetry of the action there corresponds a conserved quantity (1915AC)
- (Jose Enrique) Moyal star product on a manifold (1949AC)
- (Masud) Chaichian proof of (Hendrik) Lorentz invariance of Quantum Field Theories based on Non-Commutative Geometry(2004AC)



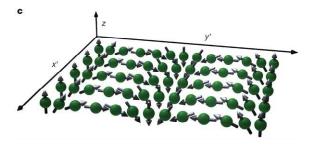






<sup>\*</sup>Fajnšmekers only

### Persistent helix



J.D.Koralek et al, Nature 458, 610 (2009).