

Superabsorption of light via quantum engineering

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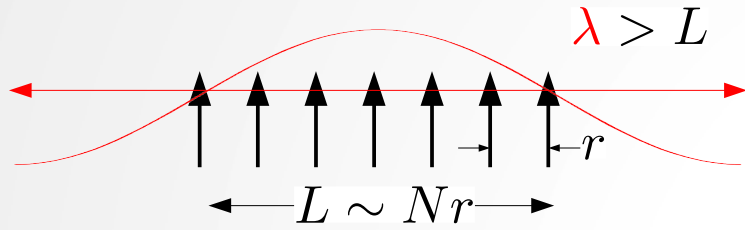
⁴*SUPA, Department of Physics, Heriot Watt University, Edinburgh EH14 4AS, United Kingdom*

ArXiv:1306.1483 (July 2013)

presented by

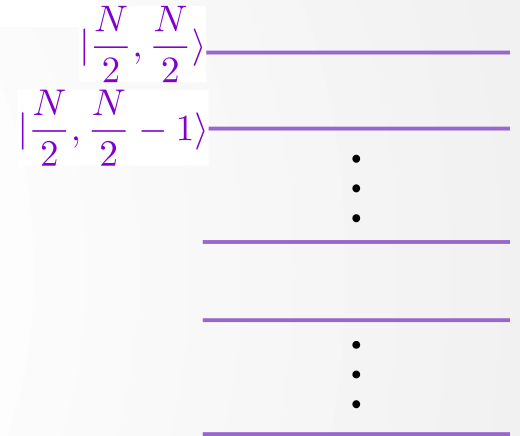
Simon E. Nigg

Superradiance in a nutshell



$$H = H_0 + q(\vec{E} \cdot \vec{d}) \sum_{j=1}^N (\sigma_j^+ + \sigma_j^-)$$

$$H_0 = \frac{\omega_A}{2} \sum_{j=1}^N \sigma_j^z$$



$$J^z = \frac{1}{2} \sum_{i=1}^N \sigma_i^z$$

$$J^\pm = \frac{1}{2} \sum_{i=1}^N \sigma_i^\pm$$

$$J^x = (J^+ + J^-)/2 \quad [J^2, J^\pm]_- = 0$$

$$J^y = -i(J^+ - J^-)/2 \quad [J^z, J^\pm]_- \neq 0$$

Dicke states: $|J, M\rangle$ ($M = -J, -J + 1, \dots, J - 1, J$)

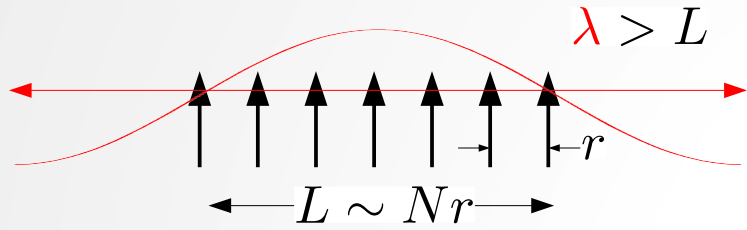
$$J^z |J, M\rangle = M |J, M\rangle$$

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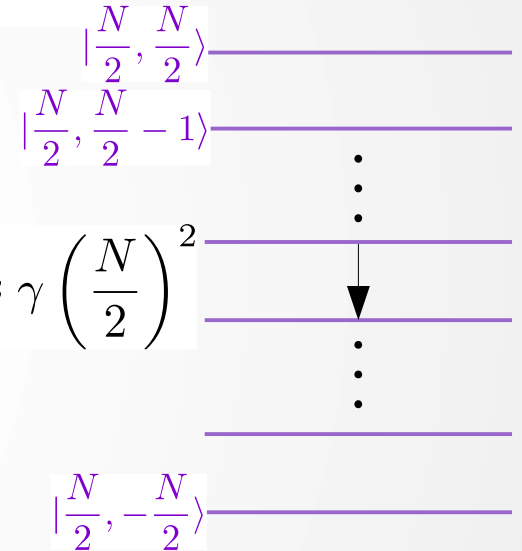
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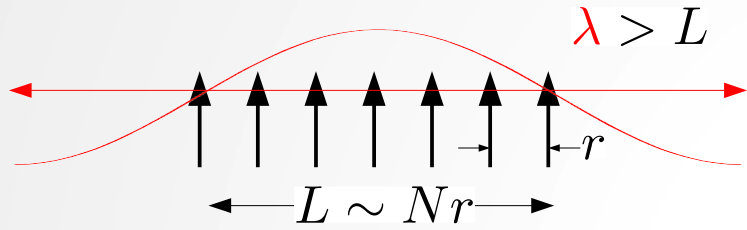
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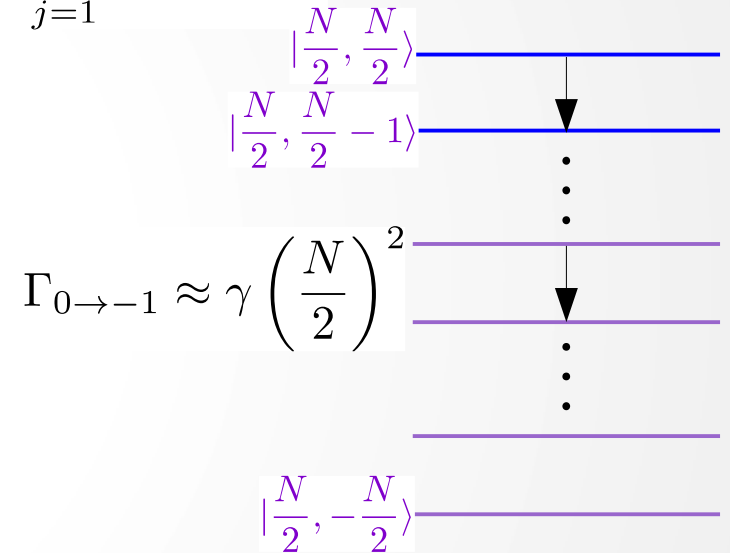
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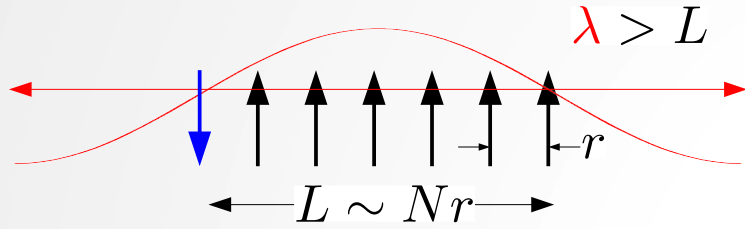
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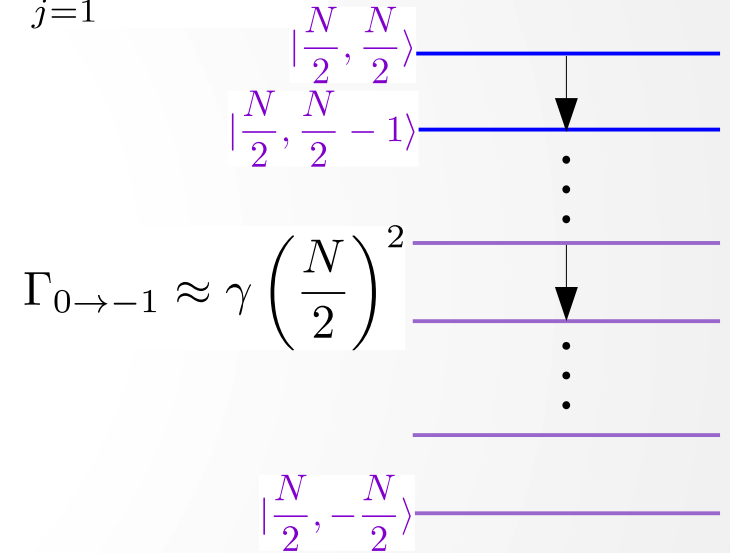
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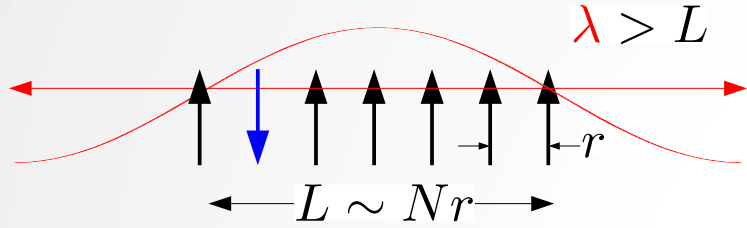
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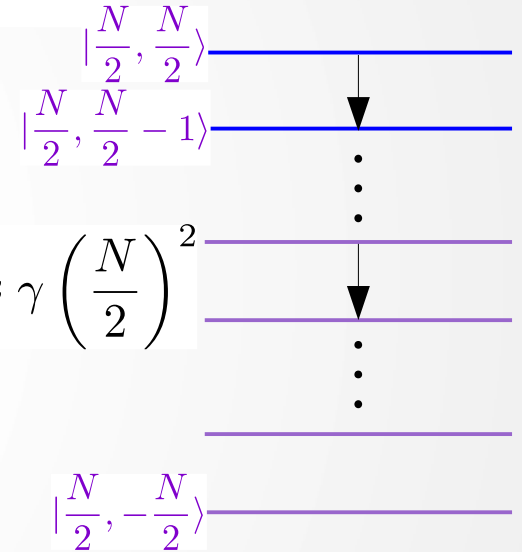
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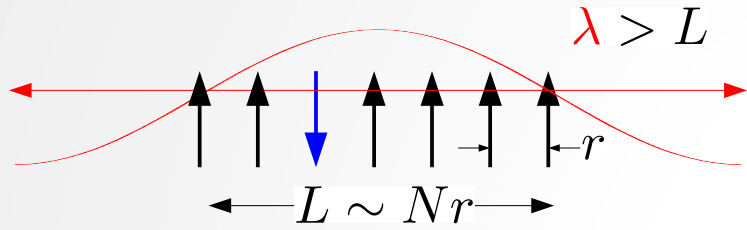
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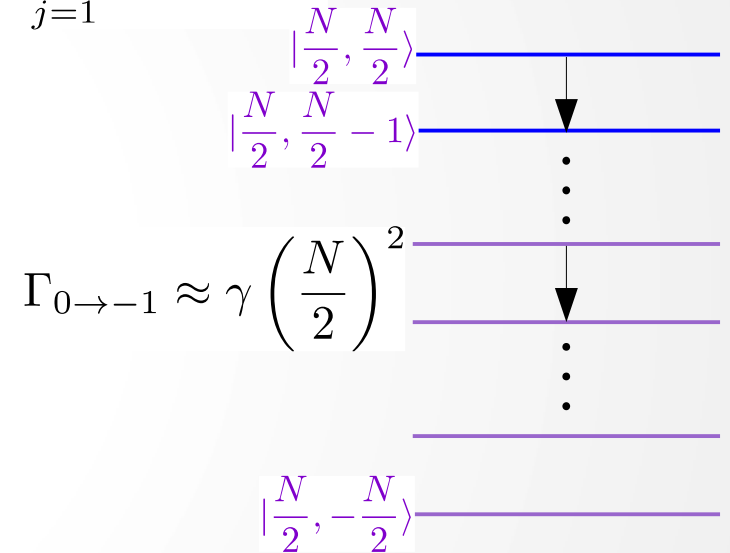
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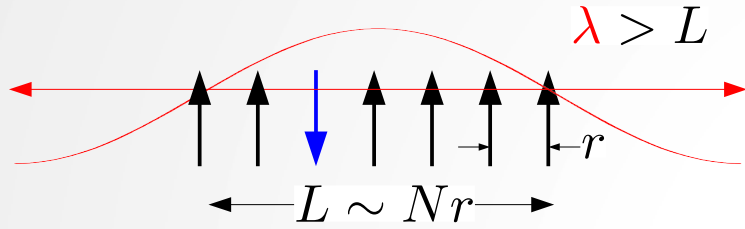
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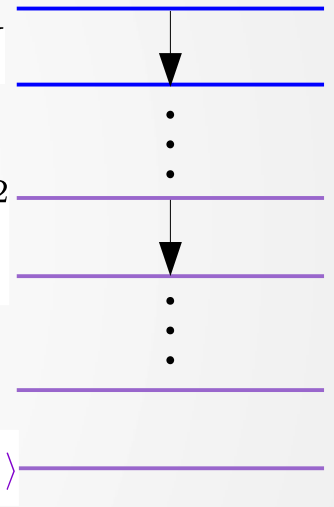
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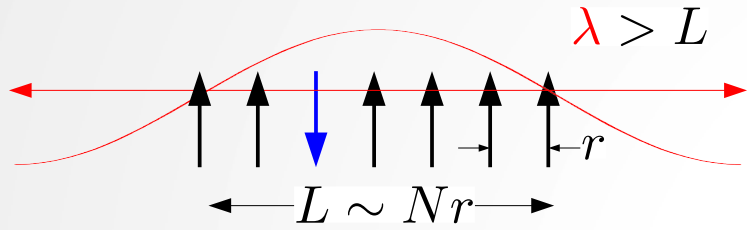
$$H_0 = \frac{\omega_A}{2} \sum_{j=1}^N \sigma_j^z \quad \Gamma_{\frac{N}{2} \rightarrow \frac{N}{2} - 1} = \gamma N$$

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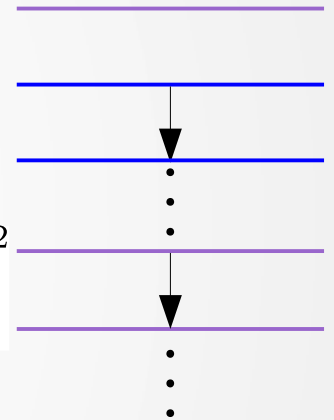
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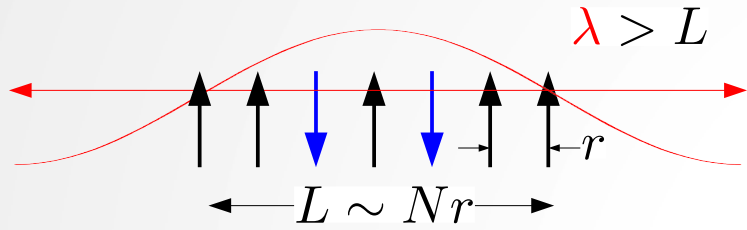
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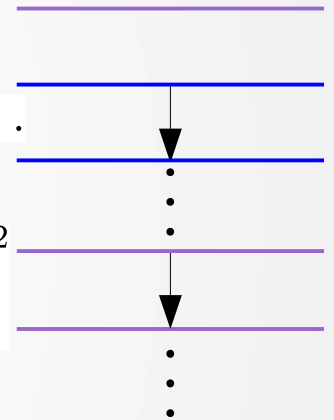
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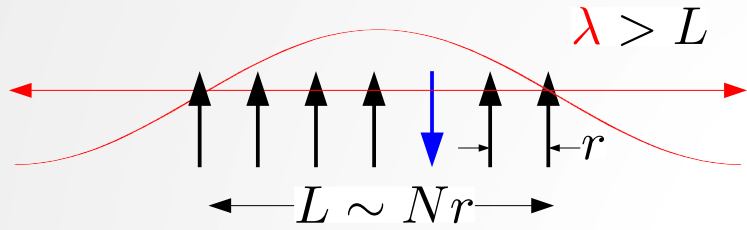
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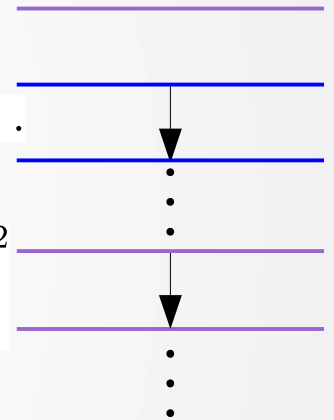
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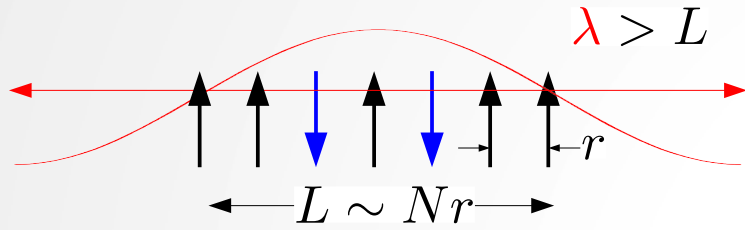
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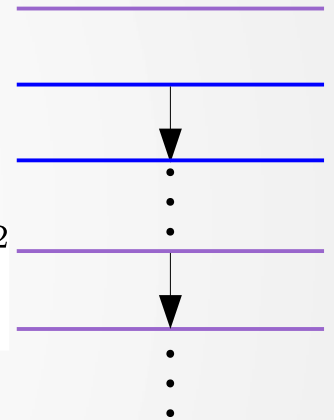
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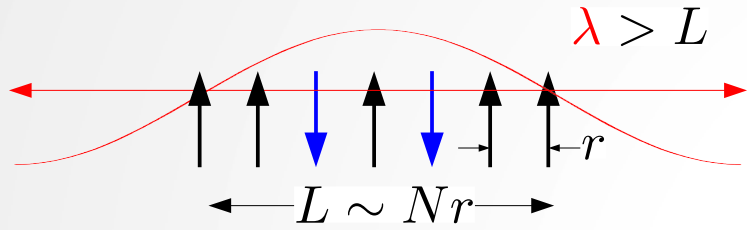
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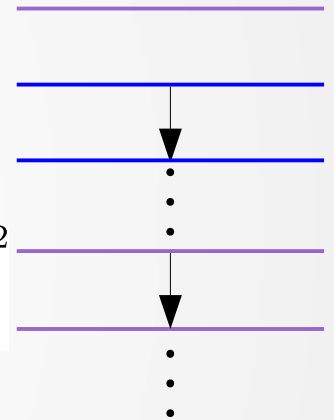
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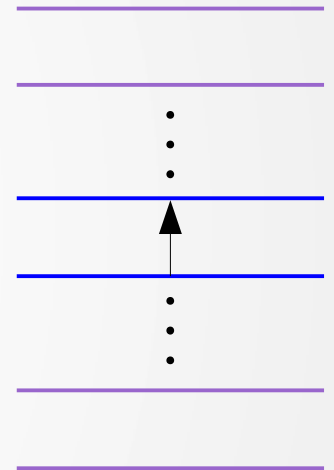


Enhanced emission rate compared to classical case for $N > 4$

From Superradiance to Superabsorption

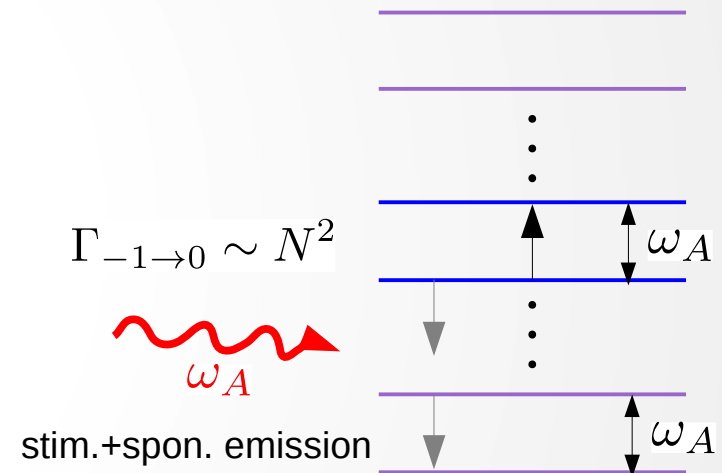
Time reversal symmetry: Superradiance \Leftrightarrow Superabsorption

$$\Gamma_{-1 \rightarrow 0} \sim N^2$$



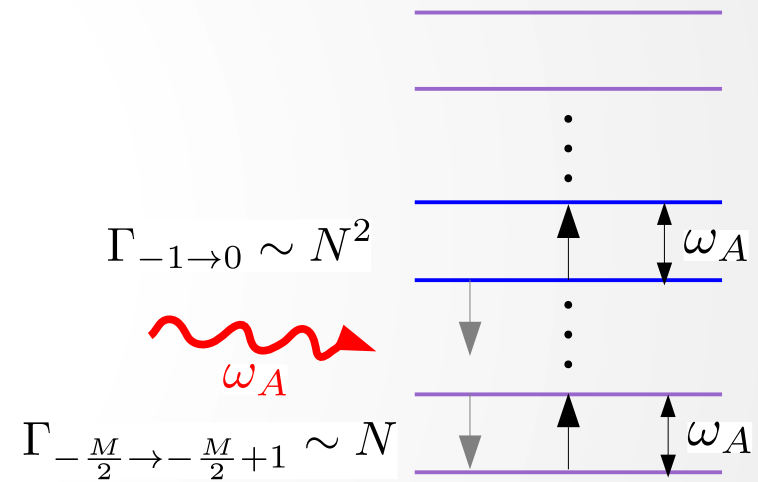
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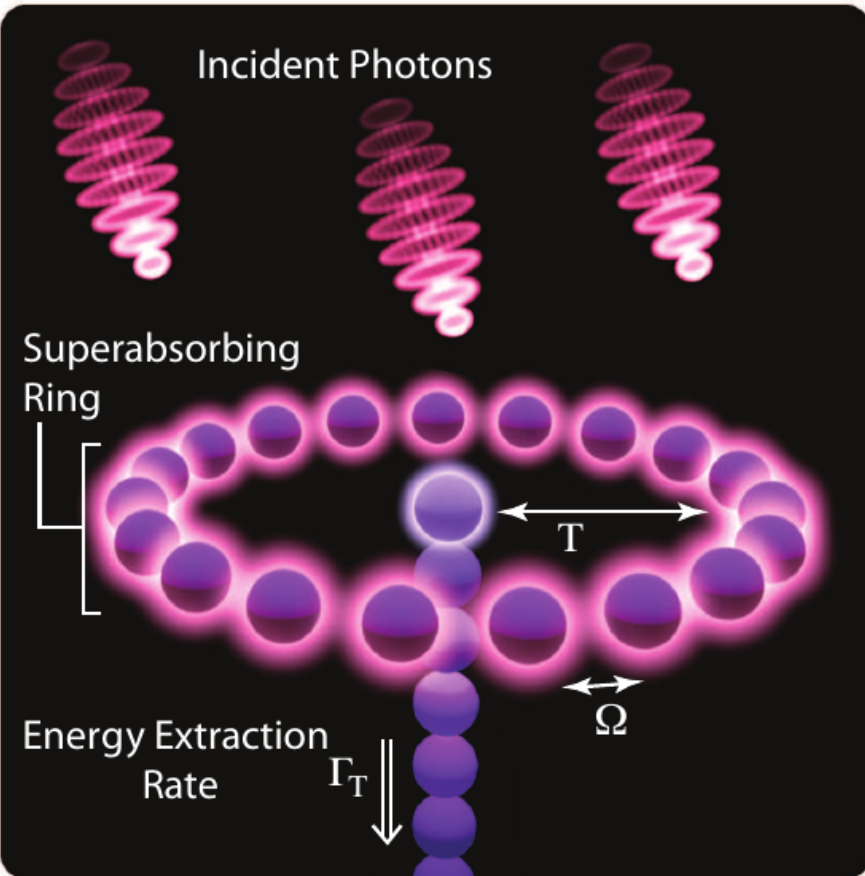


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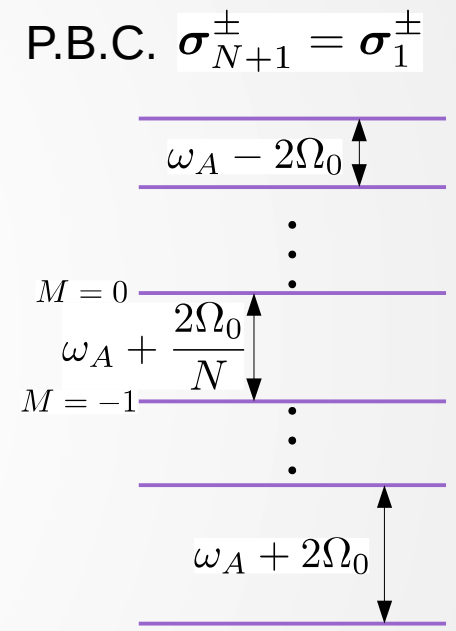


$$H_I = \Omega_0 \sum_{j=1}^N (\sigma_j^+ \sigma_{j+1}^- + h.c.)$$

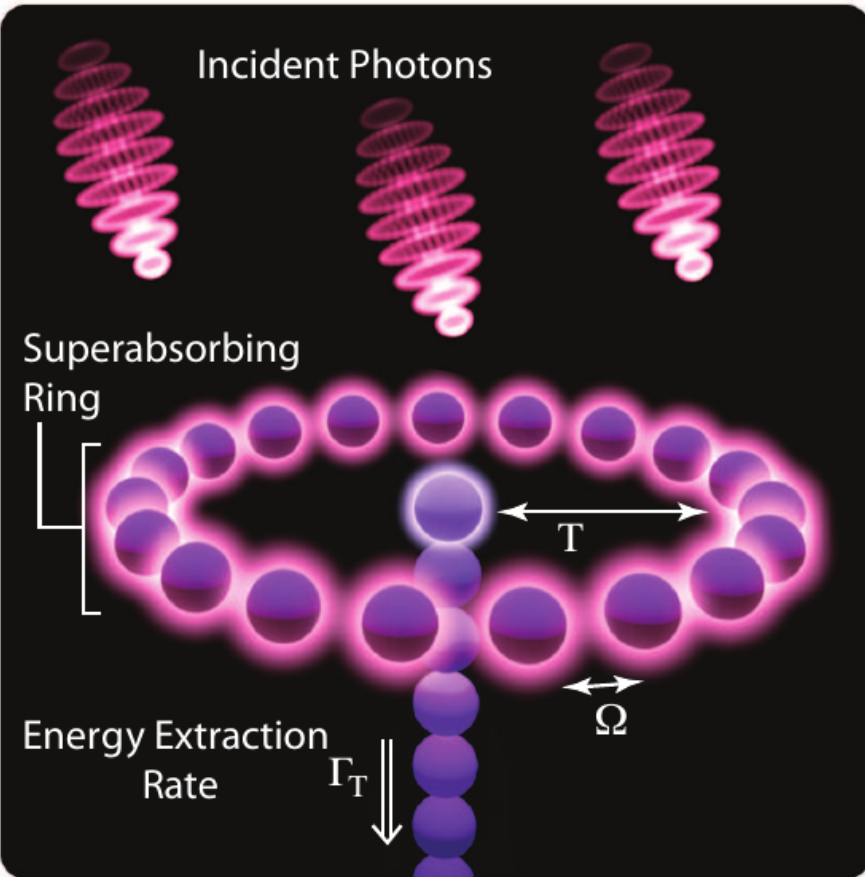
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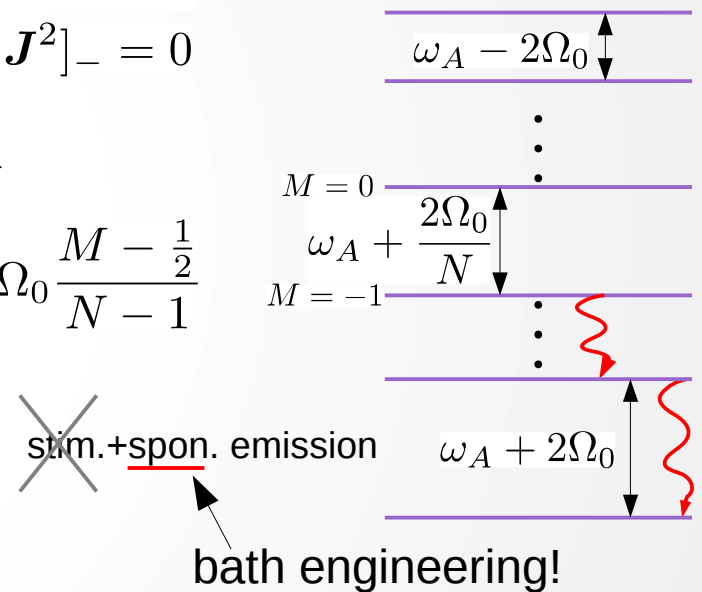
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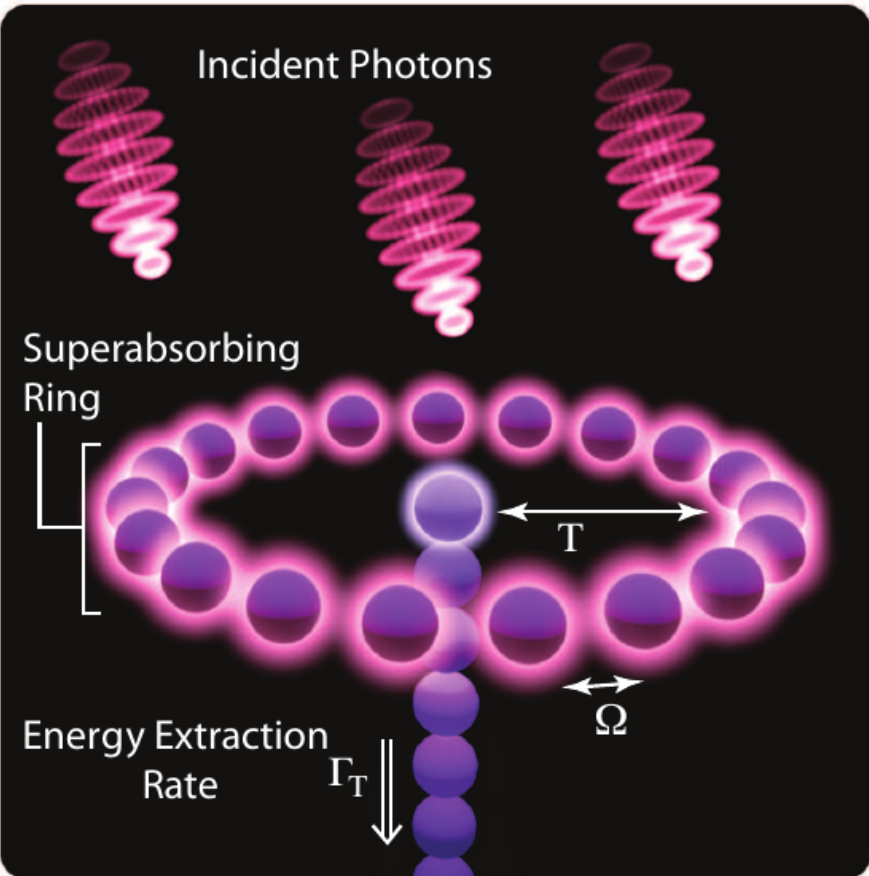
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P.B.C. $\sigma_{N+1}^\pm = \sigma_1^\pm$



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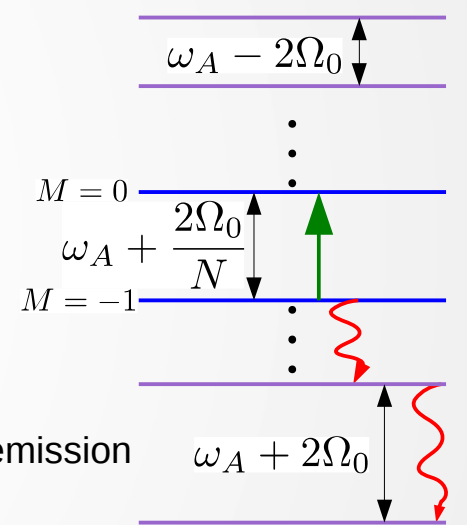
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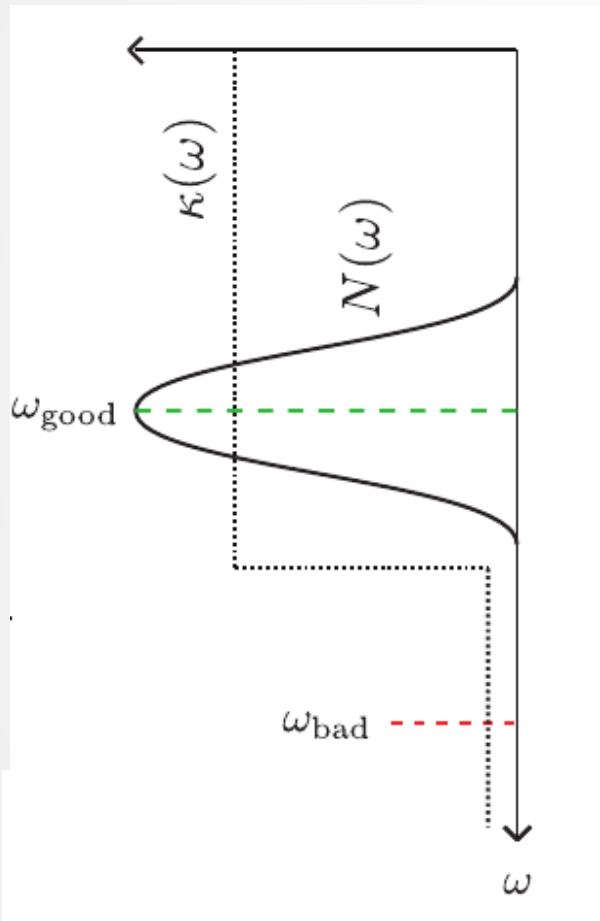


~~stim. + spon. emission~~
 bath engineering!

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$$- \gamma \sum_{\beta} \kappa(\omega_{\beta}) \left([n(\omega_{\beta}) + 1] D[\mathbf{L}_{\beta}] \rho + n(\omega_{\beta}) D[\mathbf{L}_{\beta}^{\dagger}] \rho \right)$$

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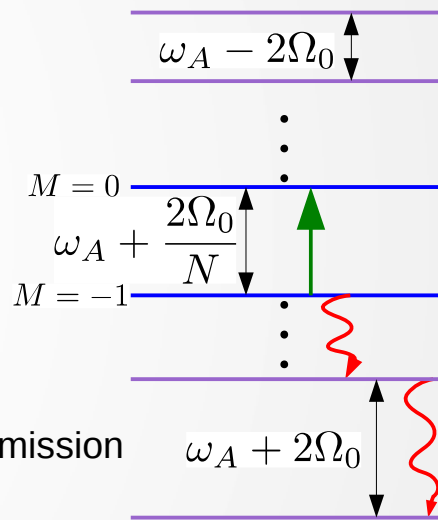
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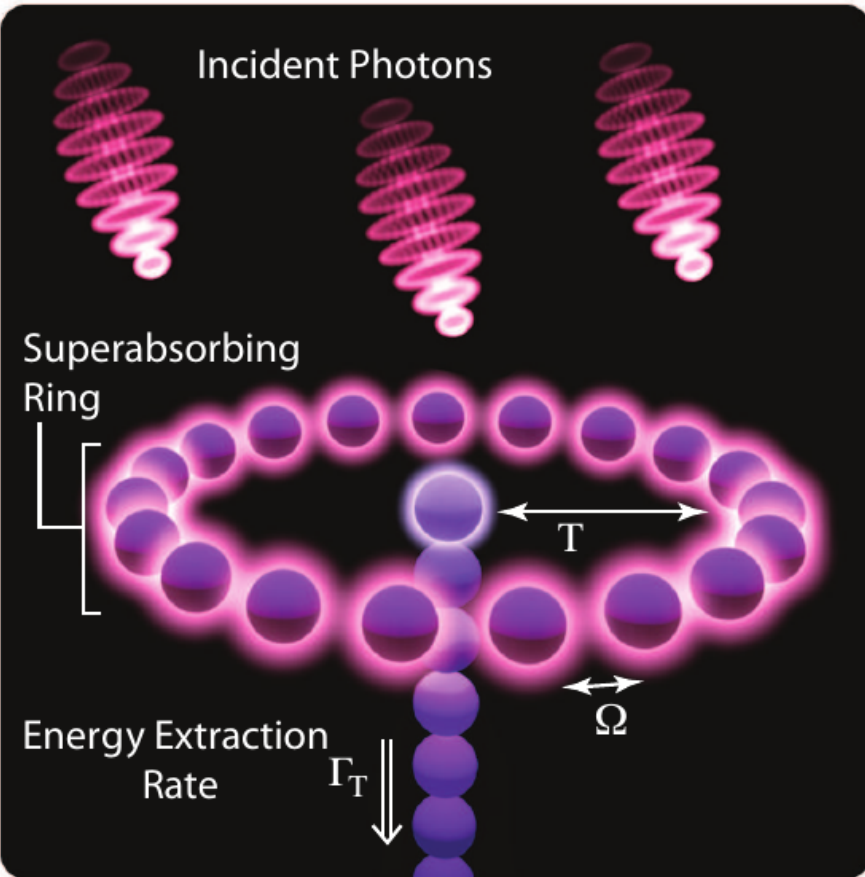
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$$\dot{\rho} = -i[\mathbf{H} + \mathbf{H}_I, \rho]$$

$$- \gamma \sum_{\beta} \kappa(\omega_{\beta}) \left([n(\omega_{\beta}) + 1] D[\mathbf{L}_{\beta}] \rho + n(\omega_{\beta}) D[\mathbf{L}_{\beta}^{\dagger}] \rho \right)$$

Towards applying superabsorption

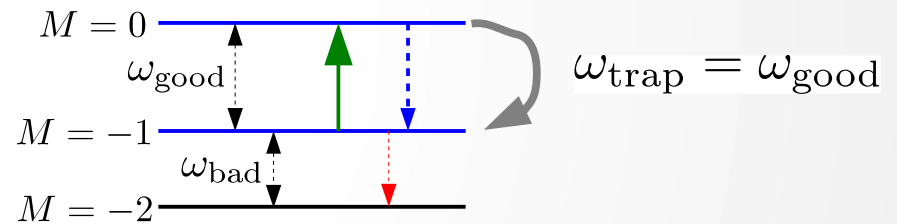


$$\mathbf{H}_T = g(\mathbf{J}^+ \boldsymbol{\sigma}_T^- + h.c.) + \omega_{\text{trap}} \boldsymbol{\sigma}_T^+ \boldsymbol{\sigma}_T^-$$

Phenomenological dissipator

$$D[\mathbf{L}_{\text{trap}}] \quad \text{with} \quad \mathbf{L}_{\text{trap}} = \sqrt{\Gamma_{\text{trap}}} \left| \frac{N}{2}, -1 \right\rangle \left\langle \frac{N}{2}, 0 \right|$$

$$\Rightarrow I_{\text{trap}}(t) = \Gamma_{\text{trap}} \text{tr} \left[\left| \frac{N}{2}, 0 \right\rangle \left\langle \frac{N}{2}, 0 \right| \rho(t) \right]$$



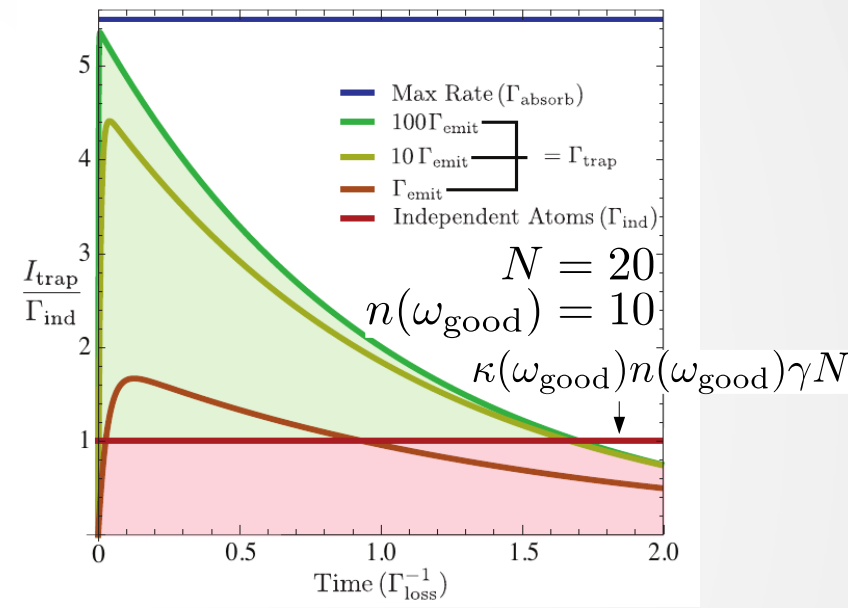
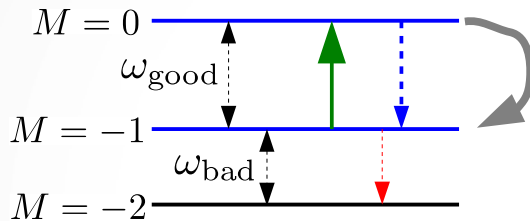
Ideal result:

$$\Gamma_{\text{loss}} = \kappa(\omega_{\text{bad}}) [n(\omega_{\text{bad}}) + 1] \Gamma_{-1 \rightarrow -2} = 0$$

$$\Gamma_{\text{trap}} \gg \Gamma_{\text{emit}} = \kappa(\omega_{\text{good}}) [n(\omega_{\text{good}}) + 1] \Gamma_{0 \rightarrow -1}$$

$$\Rightarrow I_{\text{trap}} = \Gamma_{\text{absorb}} \stackrel{N \gg 1}{\approx} \gamma \kappa(\omega_{\text{good}}) n(\omega_{\text{good}}) \left(\frac{N}{2} \right)^2$$

Results for finite Γ_{loss}



1. Assuming ideal optimal feedback

$$I_{\text{trap}} = \Gamma_{\text{absorb}} - \Gamma_{\text{loss}} = (\mu - \sigma) \left(\frac{N}{2}\right)^2 + 2\sigma$$

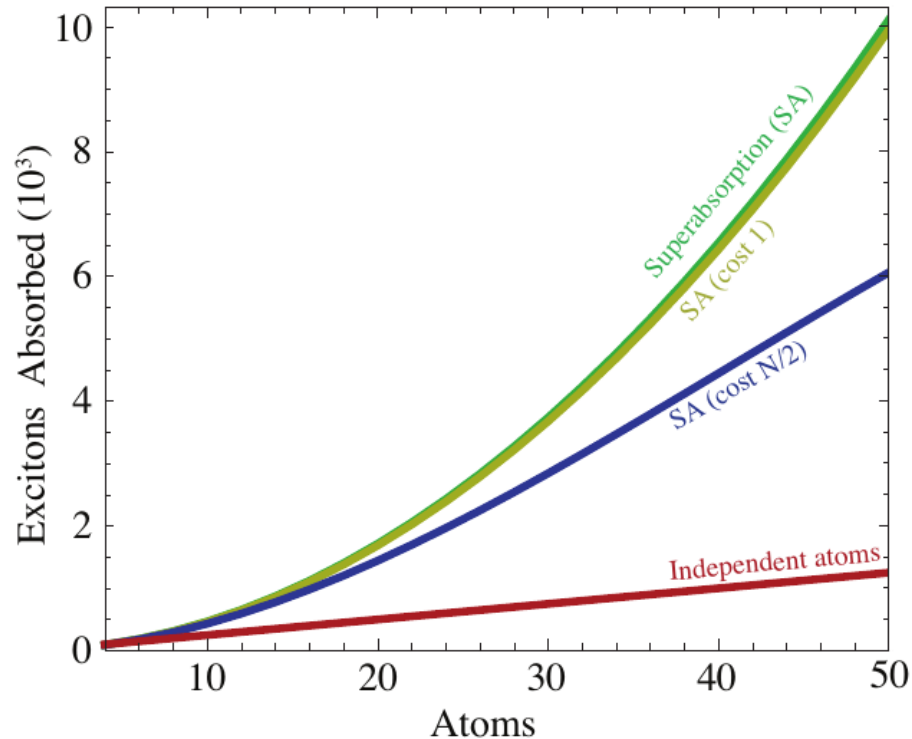
$$\mu = \gamma \kappa(\omega_{\text{good}}) n(\omega_{\text{good}})$$

$$\sigma = \gamma \kappa(\omega_{\text{bad}}) [n(\omega_{\text{bad}}) + 1]$$

\Rightarrow Superabsorption for $\mu > \sigma$

Results for finite Γ_{loss}

2. Reinitialization after $\Gamma_{\text{loss}}^{-1}$



Potential systems and applications

- Molecular rings
- Quantum dots
- Circuit QED
- BEC, optical lattices

- Light harvesting (solar cells)
- Ultra-sensitive optical or microwave sensors