

# Absence of Quantum Time Crystals in Ground States

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In analogy with crystalline solids around us, Wilczek recently proposed the idea of "time crystals" as phases that spontaneously break the continuous time translation into a discrete subgroup. The proposal stimulated further studies and vigorous debates whether it can be realized in a physical system. However, a precise definition of the time crystal is needed to resolve the issue. **Here we first present a definition of time crystals based on the time-dependent correlation functions of the order parameter. We then prove a no-go theorem that rules out the possibility of time crystals defined as such, in the ground state of a general Hamiltonian which consists of only short-range interactions.**

arxiv:1410.2143

# Wilczek's proposal

Particle on a ring, threaded by a flux:

Lagrangian:  $L = \frac{1}{2}\dot{\phi}^2 + \alpha\dot{\phi}$

Hamiltonian:  $H = \frac{1}{2}(\pi_\phi - \alpha)^2$

Eigenstates:  $|\ell\rangle = e^{i\ell\phi}$  ( $\ell$  integer)

$$\langle \ell | H | \ell \rangle = \frac{1}{2}(\ell - \alpha)^2$$

Ground state:  $\langle \ell | \dot{\phi} | \ell \rangle = \ell - \alpha \neq 0$  (for non-integer  $\alpha$ )

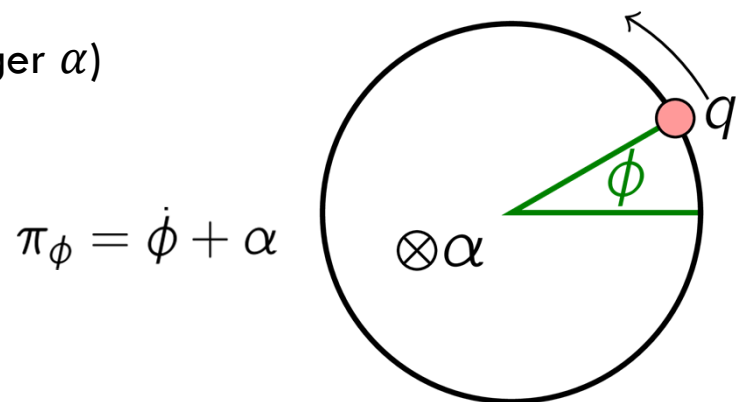
Contradiction to Heisenberg equation?

**No time-dependent observables:**

- The particle is delocalized
- The current is constant over time
- **How to get localized particles?**

Symmetry and its spontaneous breaking is a central theme in modern physics. Perhaps no symmetry is more fundamental than time-translation symmetry, since time-translation symmetry underlies both the reproducibility of experience and, within the standard dynamical frameworks, the conservation of energy. So it is natural to consider the question, whether time-translation symmetry might be spontaneously broken in a closed quantum-mechanical system. That is the question we will consider, and answer affirmatively, here.

F. Wilczek, PRL **109**, 160401 (2012)



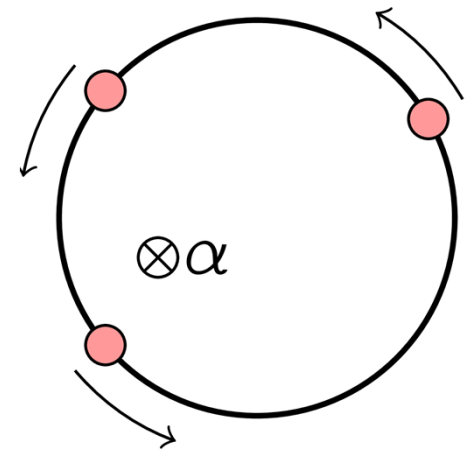
$$\pi_\phi = \dot{\phi} + \alpha$$

discussed in Fabio's JC talk 05/02/2013

# The proposed example

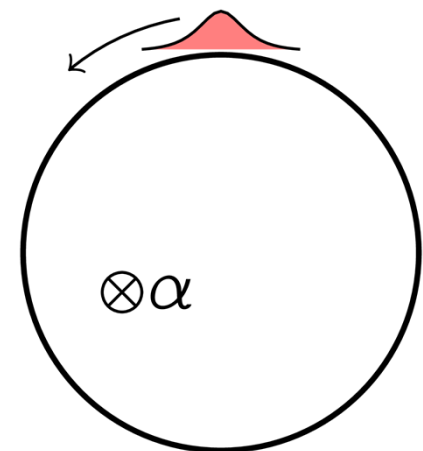
$N$  particles with weak attractive interactions:

$$H = \sum_{j=1}^N \frac{1}{2} (\pi_j - \alpha)^2 - \frac{\lambda}{N-1} \sum_{j \neq k} \delta(\phi_j - \phi_k)$$



## Wilczek's solution:

- First,  $\alpha = 0$
- Mean-field approximation (effective one-body potential)
- Gauge transformation to get solution for nonzero  $\alpha$
- Ground state is a **rotating soliton**



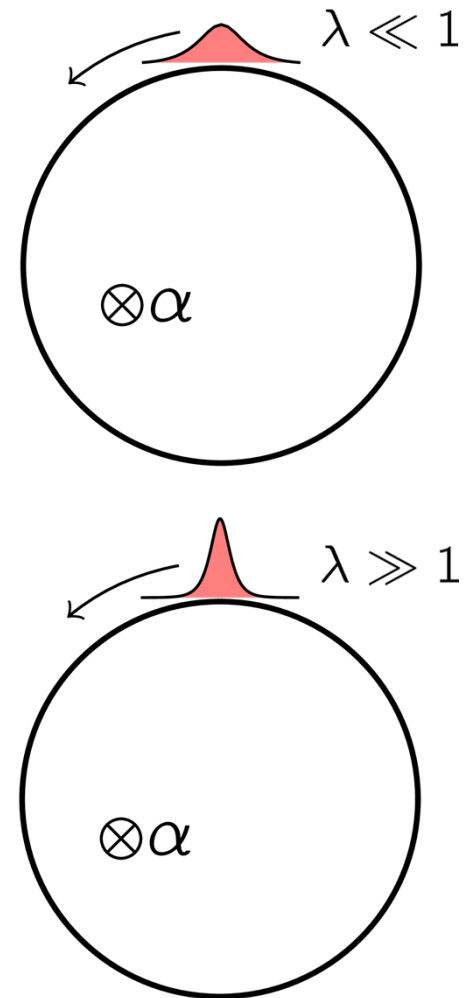
# Early criticism

## Problems with Wilczek's solution:

- In the limit  $\lambda \rightarrow \infty$ , the soliton should be insensitive to the flux
- The rotating lump should radiate, losing energy

These remarks strongly suggest that Wilczek's rotating soliton is *not* the ground state and that the true ground state is actually a stationary state, as I show below. The solution of the NLSE for arbitrary flux [3] is too lengthy and technical to fit in this Comment; thus, I shall give here

In the light of the above discussion, it seems that the very existence of "quantum time crystals" remains highly speculative.



# Spontaneous symmetry breaking

**Bogoliubov's** prescription to find spontaneous symmetry breaking in the thermodynamic limit:

- Add **symmetry-breaking perturbation** with strength  $h$
- Take the thermodynamic limit  $V \rightarrow \infty, N/V = \text{const.}$
- Take the limit  $h \rightarrow 0$

$$\lim_{h \rightarrow 0} \lim_{V \rightarrow \infty} \langle 0 | \rho(x) | 0 \rangle$$

Works great for **AFM, FM, crystalline** order, etc.

## Caveats:

- Special Hamiltonian
- Rotating frame with time-independent  $H$  exists.

## For time crystals:

Response to a **time-dependent** perturbation:

$$H = \sum_{i=1}^N \left[ \frac{(\ell_i - \alpha_i)^2}{2m_i R^2} + h_i (\phi_i - \Omega t) \right] + \sum_{j \neq k} U_{jk} (\phi_j - \phi_k)$$

Energy per particle:

$$\lim_{h \rightarrow 0} \lim_{V \rightarrow \infty} \frac{1}{N} \left( E_{\alpha,0}^{(\Omega)} - E_{\alpha,0}^{(0)} \right) > 0$$

Non-rotating **ground state** always has the smallest energy. Only for excited states, the minimum can have nonzero  $\Omega$ .

# Definition in terms of long-range order

Define spontaneous symmetry breaking using  
**long-range order of correlation functions:**

$$\sigma^2 := \lim_{|x-x'|\rightarrow\infty} \lim_{V\rightarrow\infty} \langle 0 | \phi(x, 0) \phi(x', 0) | 0 \rangle$$

**Example:** Crystalline order (periodic in space)

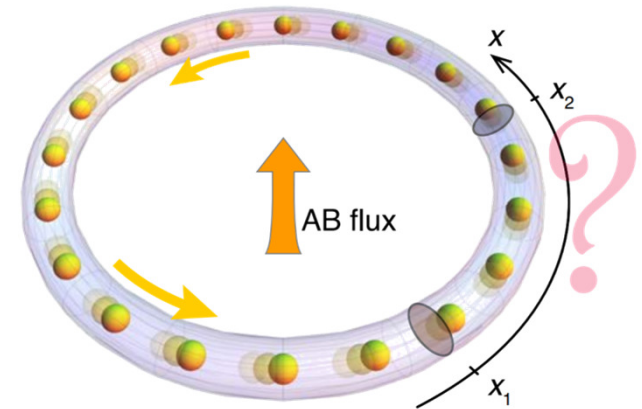
$$\lim_{|x-x'|\rightarrow\infty} \lim_{V\rightarrow\infty} \langle 0 | \rho(x, 0) \rho(x', 0) | 0 \rangle = f(x - x')$$

**By analogy, define time crystals as**

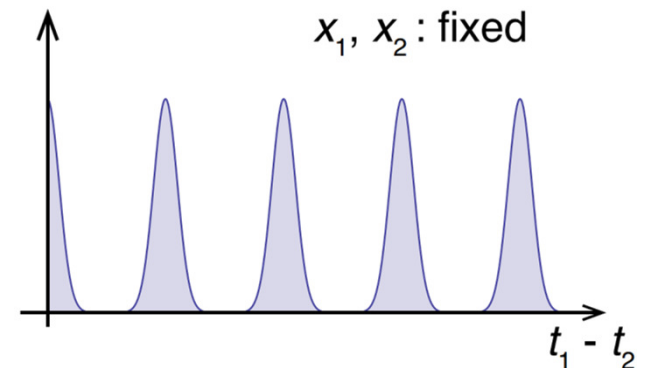
$$\lim_{|x-x'|\rightarrow\infty} \lim_{V\rightarrow\infty} \langle 0 | \phi(x, t) \phi(x', t') | 0 \rangle = f(t - t')$$

Why long-range **spatial** order?

➤ Otherwise too many “trivial” time crystals exist.



$$\langle \rho(x_1, t_1) \rho(x_2, t_2) \rangle$$



Watanabe and Oshikawa,  
arxiv:1410.2143

# Does (such) a time crystal exist?

## Assumptions:

- Locality of the Hamiltonian
- Locality of commutator

$$H = \int d^d x \mathcal{H}(x)$$

$$[\phi(x, t), \phi(x', t)] \approx \delta(x - x')$$

$$\lim_{|x-x'| \rightarrow \infty} \lim_{V \rightarrow \infty} \langle 0 | \phi(x, t) \phi(x', t') | 0 \rangle = f(t - t')$$

$$\lim_{V \rightarrow \infty} \frac{1}{V^2} \langle 0 | \Phi(t) \Phi(t') | 0 \rangle = f(t - t')$$

Integrated order parameter:

$$\Phi(t) = \int d^d x \phi(x, t)$$

**Step 1:** A time crystal has long-range order

$$(*) \quad \lim_{V \rightarrow \infty} \int_0^\infty d\Omega \rho(\Omega) e^{-i\Omega t} = f(t) \quad \rho(\Omega) = \frac{1}{V^2} \sum_n |\langle 0 | \Phi(0) | n \rangle|^2 \delta(\Omega - E_n) \geq 0$$

$$\implies 0 \leq |f(t)| \leq f(0)$$

$$\implies \text{Spatial long-range order because } f(0) > 0$$

# Does (such) a time crystal exist?

$$(*) \quad \lim_{V \rightarrow \infty} \int_0^\infty d\Omega \rho(\Omega) e^{-i\Omega t} = f(t)$$

**Step 2:** Investigate the state  $|\Phi\rangle = \frac{\Phi|0\rangle}{|\Phi|0\rangle|}$

$$\langle \Phi | H | \Phi \rangle = \frac{\langle 0 | [[\Phi(0), H], \Phi(0)] | 0 \rangle}{2 \langle 0 | \Phi(0)^2 | 0 \rangle} \propto \frac{V^3 V^{-2}}{f(0) V^2} = O(V^{-1})$$

$$(**) \quad \int_0^\infty d\Omega \rho(\Omega) \Omega = O(V^{-1}) \quad \rho(\Omega) = \frac{1}{V^2} \sum_n |\langle 0 | \Phi(0) | n \rangle|^2 \delta(\Omega - E_n) \geq 0$$

The results (\*) and (\*\*) are mutually contradictory. Therefore,  $f(t) \equiv 0$ , so **no time crystals exist** (at least according to the authors' definition).

- Small  $O(V^{-1})$  order parameter can exist in finite systems
- The proof remains true also for space-time crystals
- It does not hold for excited states or nonequilibrium states



# Conclusions



- Conventional concepts of spontaneous symmetry breaking are hard to generalize to time crystals.
- The authors' definition is based on the long-distance asymptotic behavior of correlation functions.
- According to this definition, local Hamiltonians cannot break time translation invariance spontaneously.