

Impurity Bound States and Greens Function Zeroes as Local Signatures of Topology

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We show that the local in-gap Greens function of a band insulator $G_0(\epsilon, \mathbf{k}_{\parallel}, \mathbf{r}_{\perp} = 0)$, with \mathbf{r}_{\perp} the position perpendicular to a codimension-1 or -2 impurity, reveals the topological nature of the phase. For a topological insulator, the eigenvalues of this Greens function attain zeros in the gap, whereas for a trivial insulator the eigenvalues remain nonzero. This topological classification is related to the existence of in-gap bound states along codimension-1 and -2 impurities. Whereas codimension-1 impurities can be viewed as 'soft edges', the result for codimension-2 impurities is nontrivial and allows for a direct experimental measurement of the topological nature of 2d insulators.

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Silas Hoffman

Journal club

28.4.15

Codimension?

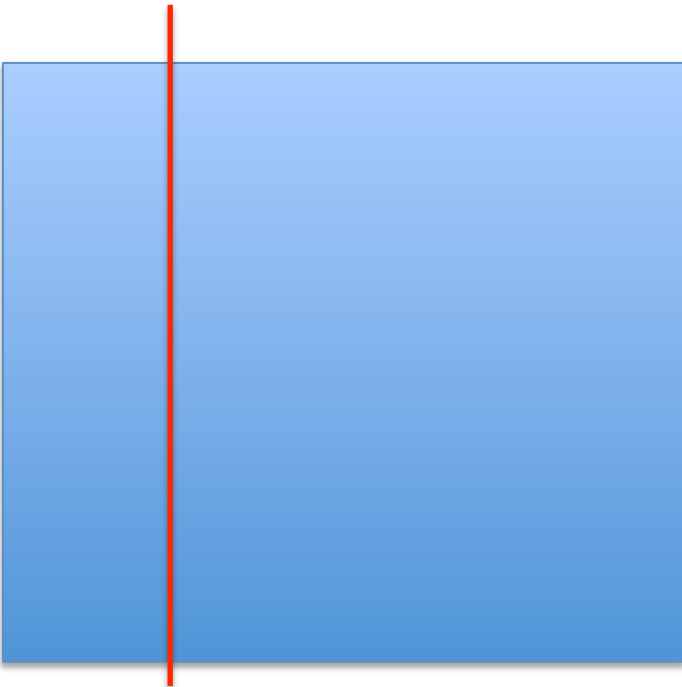
Codimension n -impurity: impurity with n fewer dimensions than the d -dimensional imbedding material

Codimension?

Codimension 1-impurity: impurity with **1** fewer dimensions than the **d-dimensional imbedding material**

2-dimensional material: $(2-1=1)$

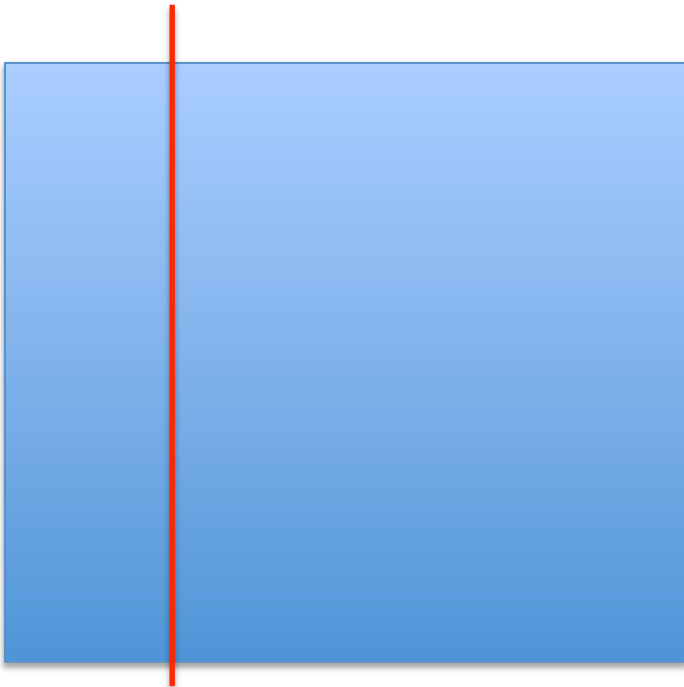
a line



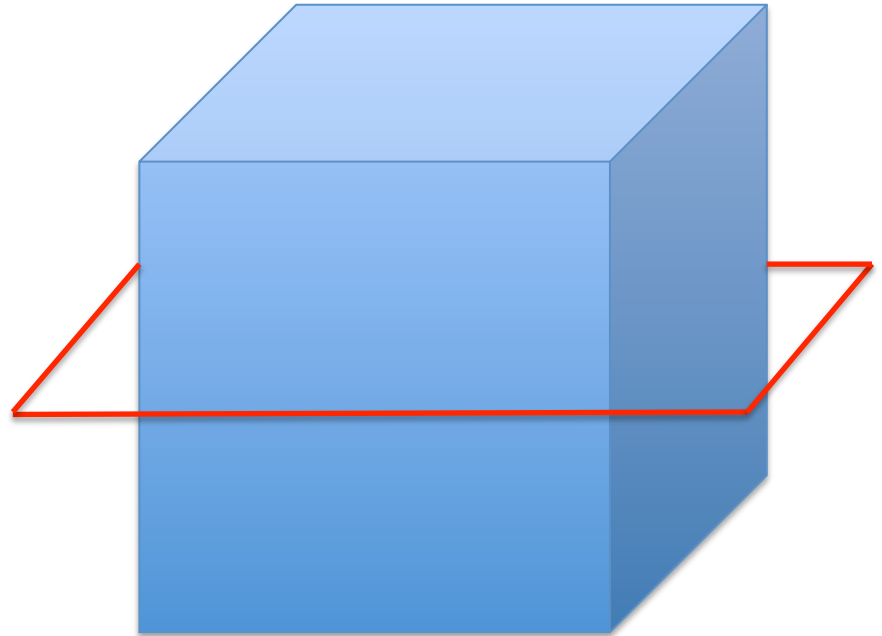
Codimension?

Codimension 1-impurity: impurity with **1 fewer dimensions** than the **d-dimensional imbedding material**

2-dimensional material: $(2-1=1)$
a line



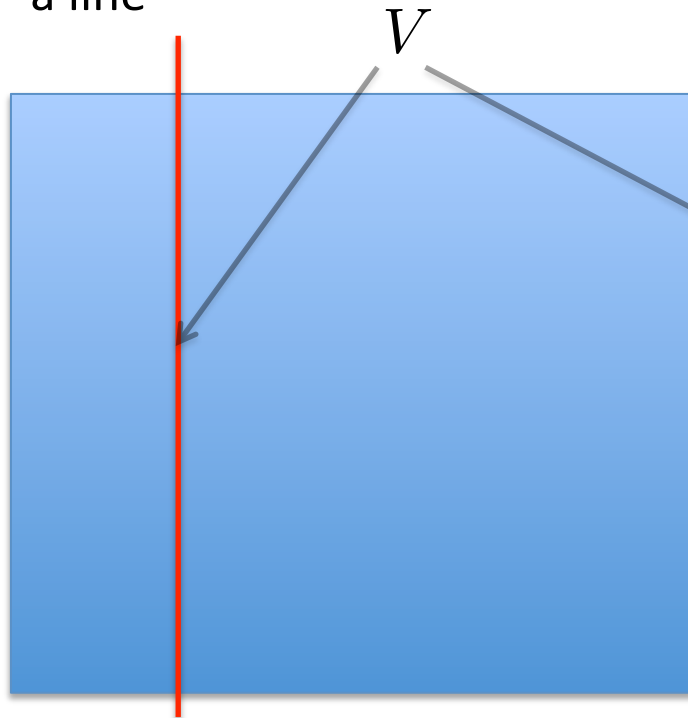
3-dimensional material: $(3-1=2)$
a plane



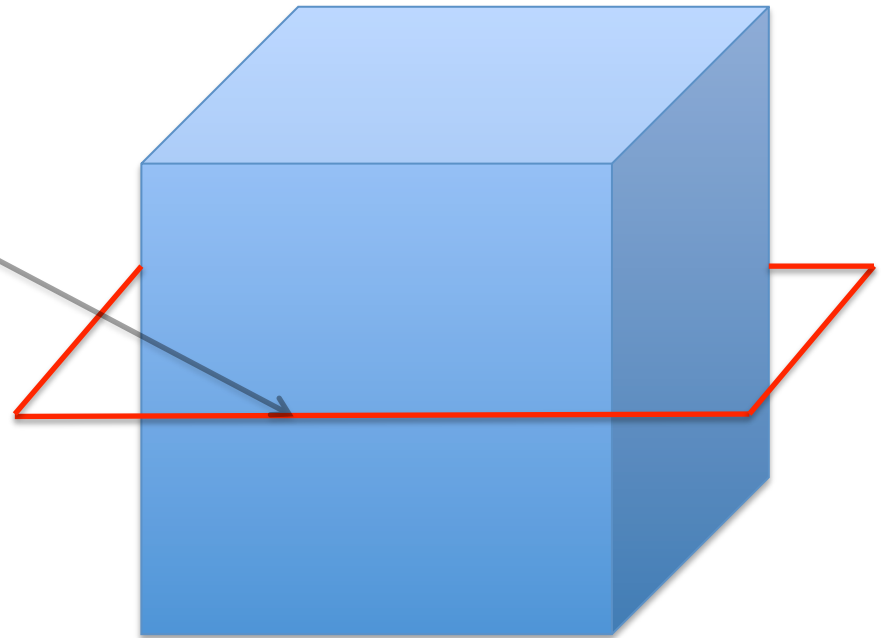
Codimension?

Codimension 1-impurity: impurity with **1 fewer dimensions** than the **d-dimensional imbedding material**

2-dimensional material: $(2-1=1)$
a line

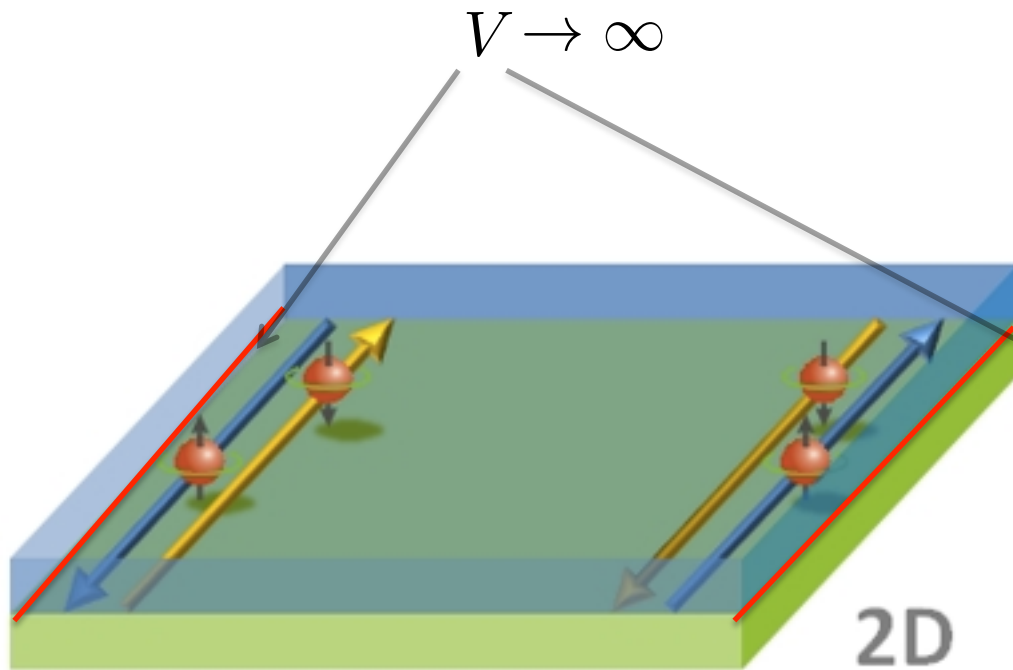


3-dimensional material: $(3-1=2)$
a plane



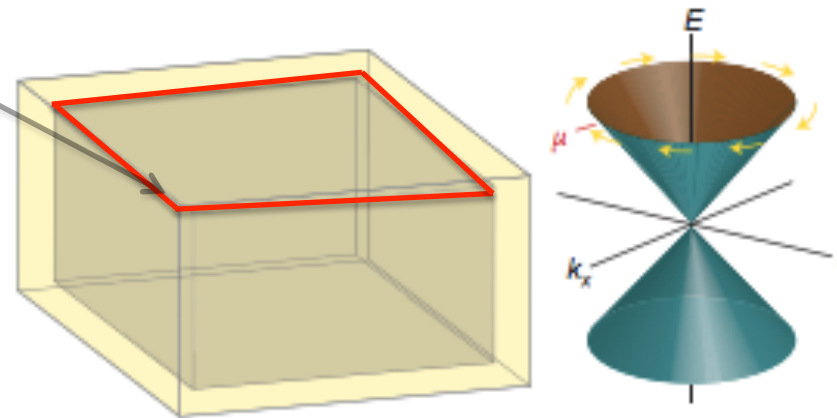
Codimension?

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When there is a ('hard') edge, we expect a localized in-gap state

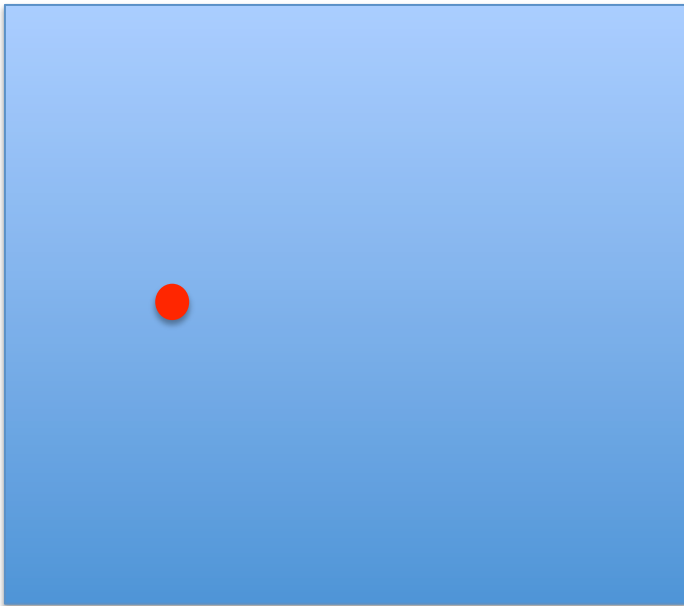
No surprise if we find a localized in-gap state at a 'soft edge'



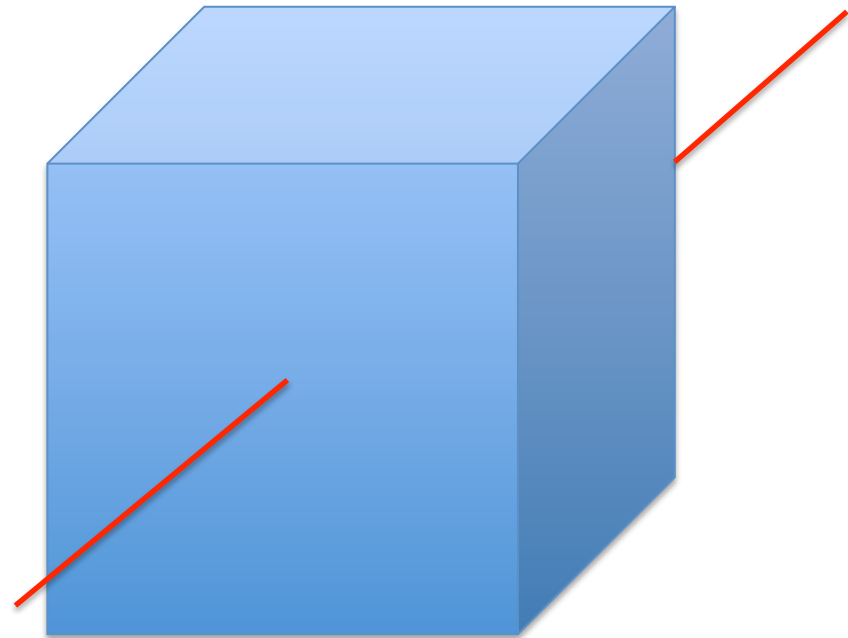
Codimension?

Codimension 2-impurity: impurity with **2 fewer dimensions** than the **d-dimensional imbedding material**

2-dimensional material: $(2-2=0)$
a point



3-dimensional material: $(3-1=2)$
a line

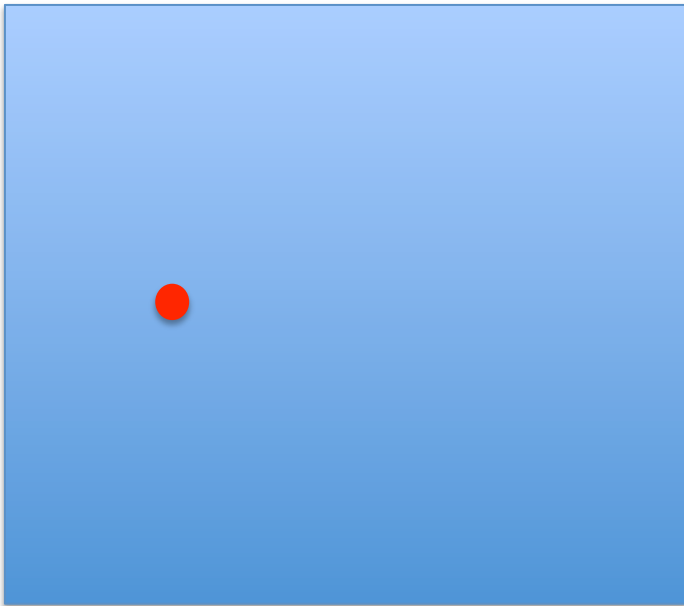


Codimension?

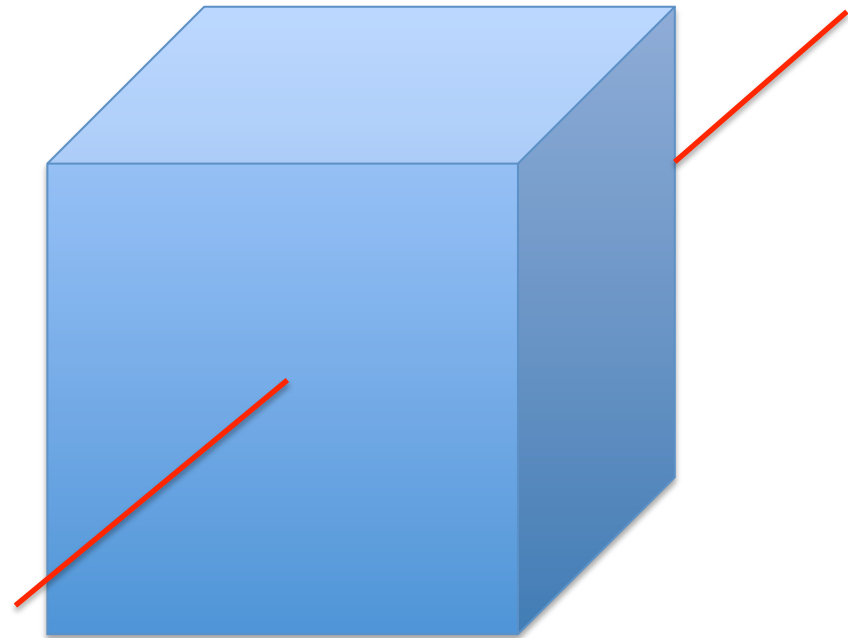
Claim: codim. 1 and 2-impurities in TI's always host bound states

Codimension 2-impurity: impurity with **2 fewer dimensions** than the **d-dimensional imbedding material**

2-dimensional material: $(2-2=0)$
a point



3-dimensional material: $(3-1=2)$
a line



Model

Theory of impurity bound states- We begin with a translationally invariant system, described by a minimal time-reversal invariant two-band model. The generic Hamiltonian assumes the form

$$H_0 = \sum_{\mathbf{k}\alpha\beta} c_{\mathbf{k}\alpha}^\dagger d_\mu(\mathbf{k}) \gamma_{\alpha\beta}^\mu c_{\mathbf{k}\beta} \quad (1)$$

$$\gamma^0 = \sigma_0 \otimes \tau_3 \text{ and } \gamma^i = \sigma_i \otimes \tau_1$$

$$d_0(\mathbf{k}) = M - 2B \sum_i (1 - \cos k_i)$$

$$d_i(\mathbf{k}) = \sin k_i$$

BHZ model: topologically nontrivial when $0 < M/B < 4d$.

Model

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$$\mathbf{G}_0(\omega, \mathbf{k}) = \frac{1}{\omega - d_\mu(\mathbf{k})\gamma^\mu} = \frac{\omega \mathbf{1}_4 + d_\mu(\mathbf{k})\gamma^\mu}{\omega^2 - |d(\mathbf{k})|^2}$$

$$H_V = \sum_{\mathbf{r}\alpha\beta} c_{\mathbf{r}\alpha}^\dagger V_{\alpha\beta}(\mathbf{r}) c_{\mathbf{r}\beta} \quad (H_0 + H_V)\psi_\epsilon(\mathbf{r}) = \epsilon\psi_\epsilon(\mathbf{r})$$

$$\psi_\epsilon(\mathbf{r}) = \sum_{\mathbf{r}'} \mathbf{G}_0(\epsilon, \mathbf{r} - \mathbf{r}') V(\mathbf{r}') \psi_\epsilon(\mathbf{r}')$$

Model

Let us therefore consider the simplest possible choice: a constant $V(\mathbf{r})$ along a n -dimensional plane in a d -dimensional system (hence codimension $d - n$). The d -dimensional position vector \mathbf{r} can be split into the perpendicular coordinates \mathbf{r}_\perp and the parallel coordinates \mathbf{r}_\parallel

$$\psi_\epsilon(\mathbf{r}) = \sum_{\mathbf{r}'} \mathbf{G}_0(\epsilon, \mathbf{r} - \mathbf{r}') V(\mathbf{r}') \psi_\epsilon(\mathbf{r}')$$

$$V(\mathbf{r}) = \mathbf{V}_0 \delta_{\mathbf{r}_\perp=0}^n$$

$$\mathbf{V}_0 = V \mathbf{1} + V_0 \gamma^0$$

According to **parity** and **time reversal (TR)** invariance

$$\det [\mathbf{G}_0(\epsilon, \mathbf{k}_\parallel, \mathbf{r}_\perp = 0) \mathbf{V}_0 - \mathbf{1}] = 0$$

At least one in-gap state, for any potential, if there is

Existence of in-gap states related to **eigenvalue**

a zero EV

Codimension-1

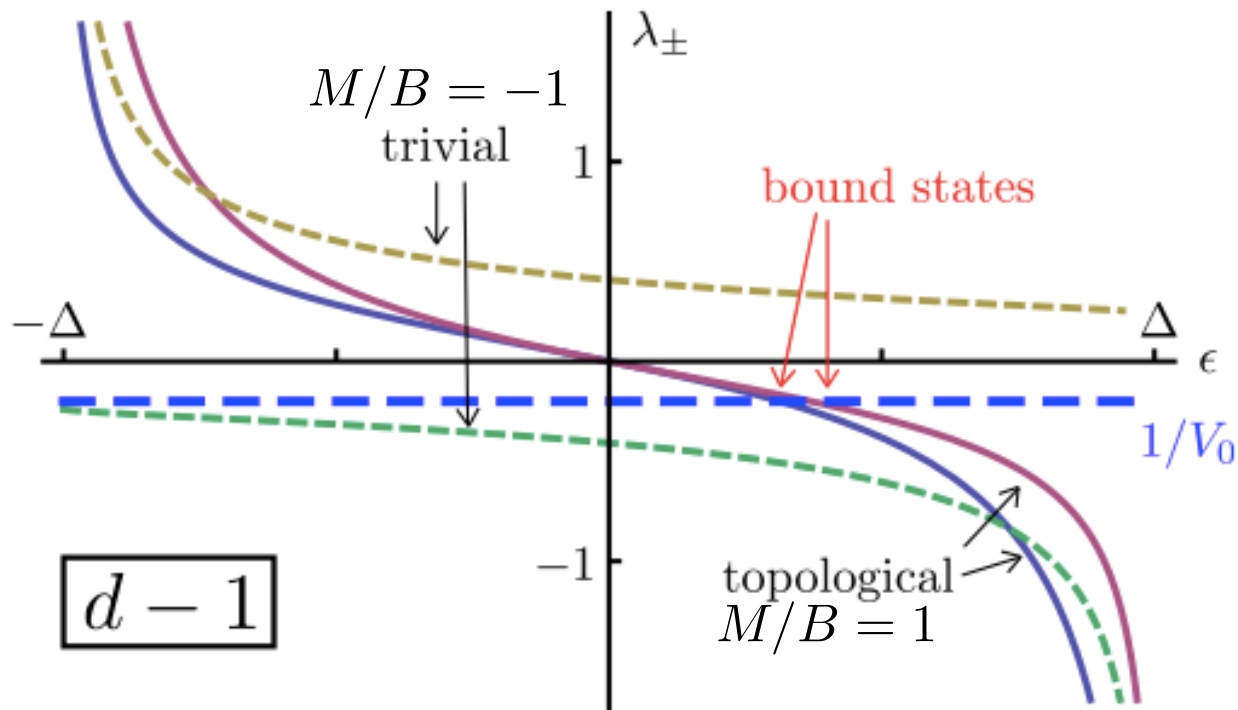
One can show:

- Normal insulator has **no zero eigenenergies**
- TI **must have a zero eigenenergies** and therefore supports at least a single **bound state**
- TI eigenvalues go to +/- infinity at the valence/cond. band

Codimension-1

One can show:

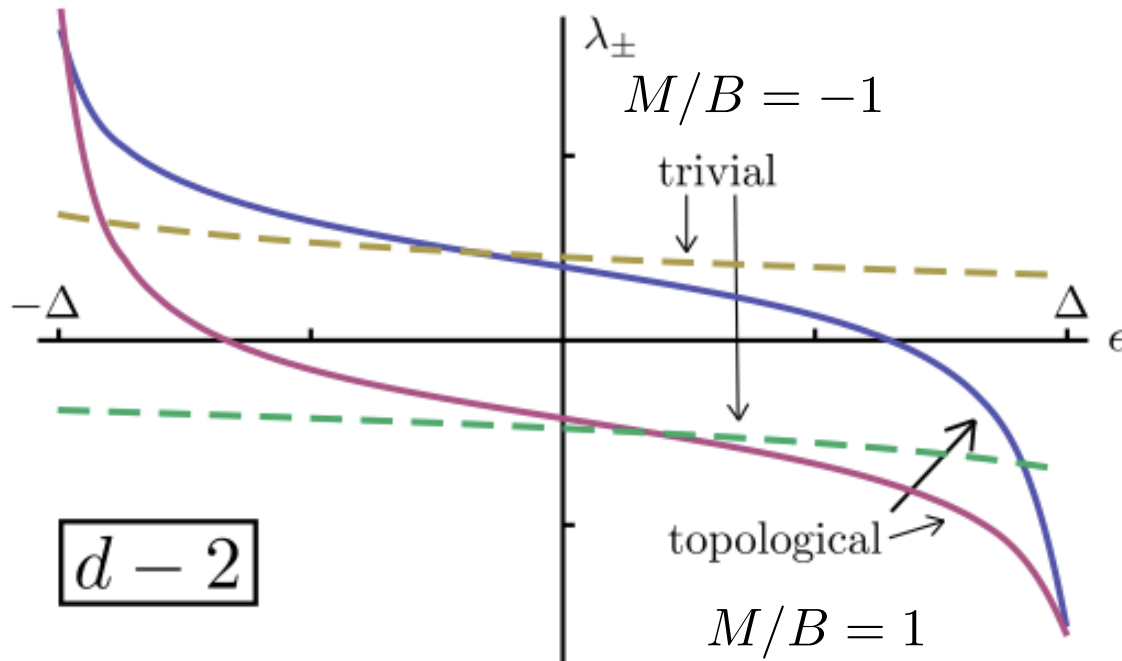
- Normal insulator has **no zero eigenenergies**
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- TI eigenvalues go to \pm infinity at the valence/cond. band



Codimension-2

One can show:

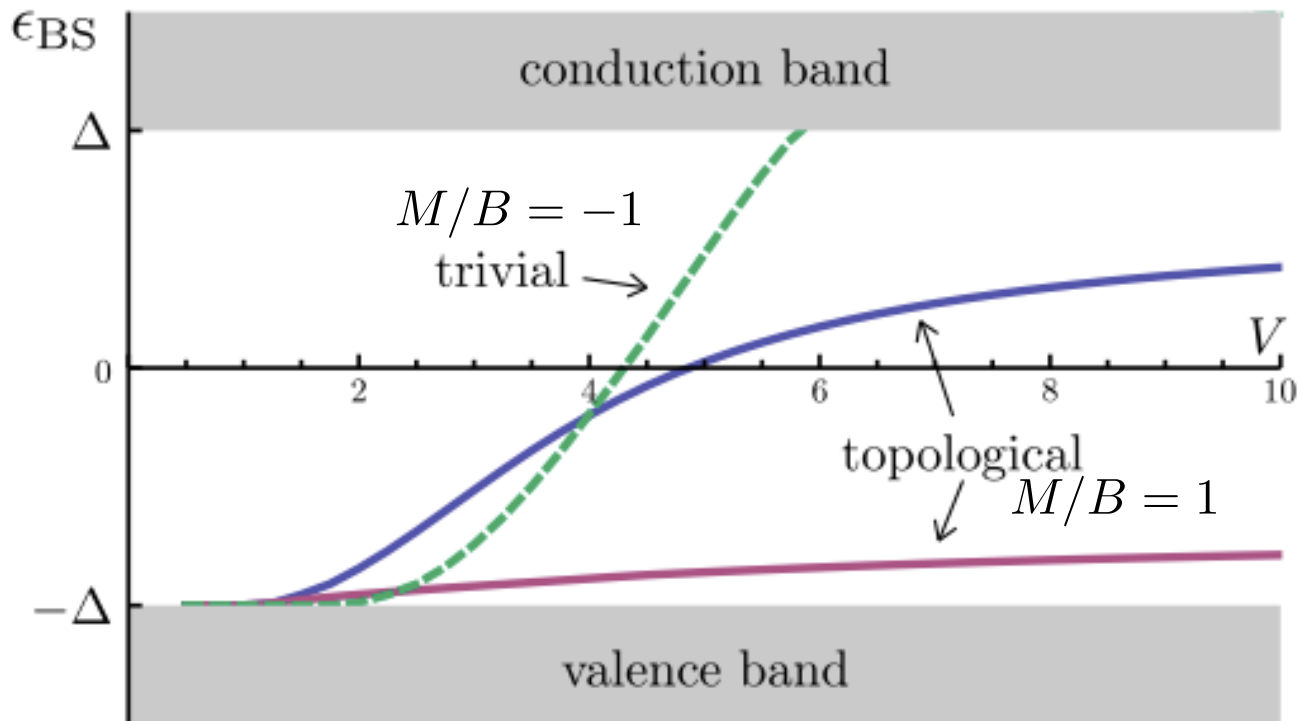
- Normal insulator has **no zero eigenenergies**
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Codimension-1

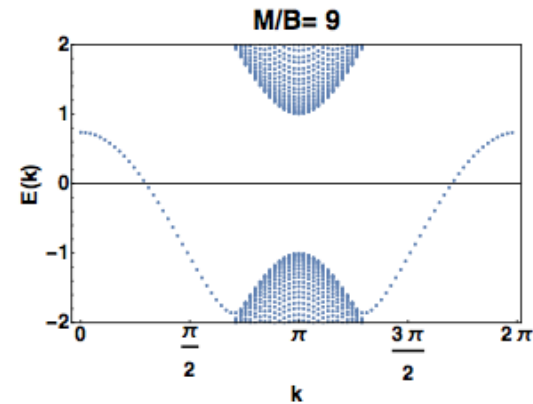
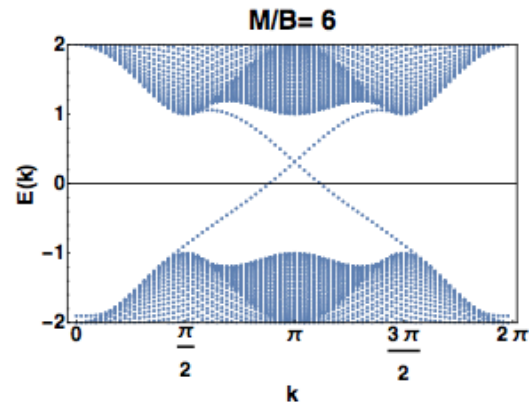
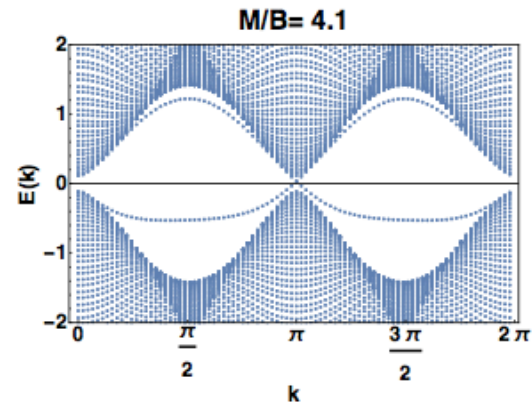
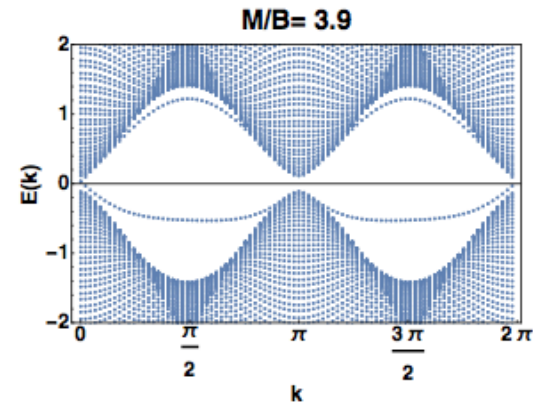
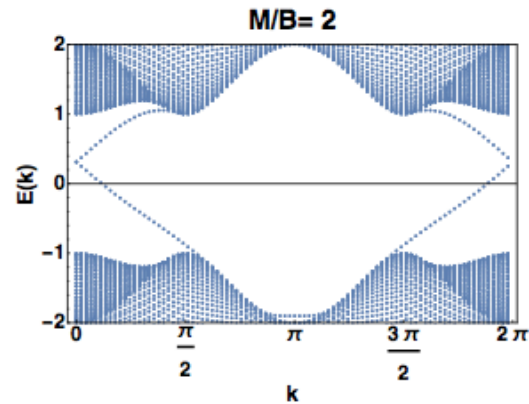
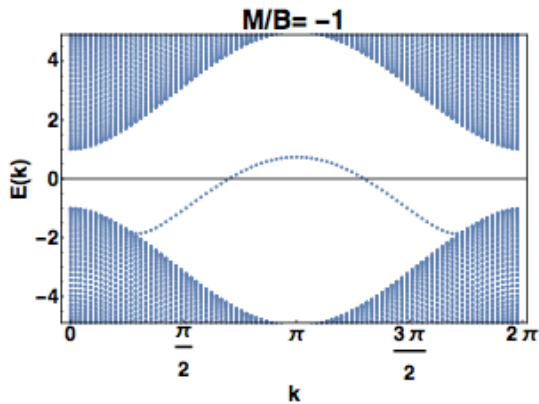
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Tight Binding (60 x 90 array) Line Impurity

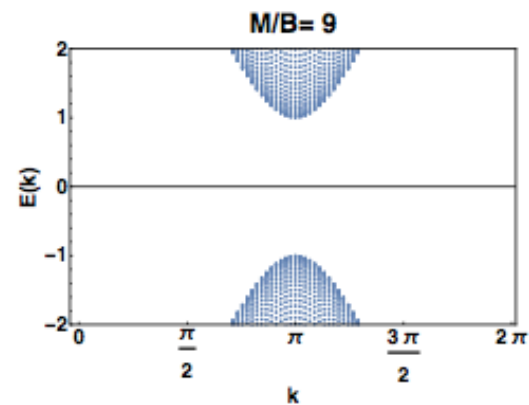
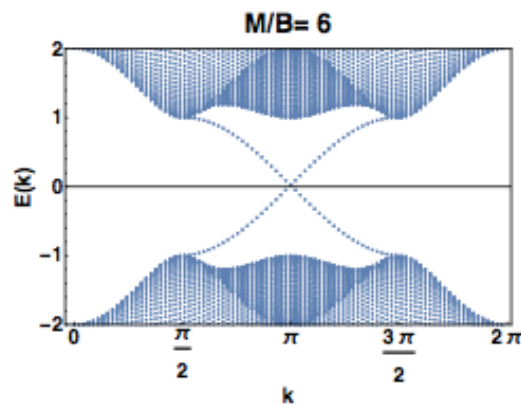
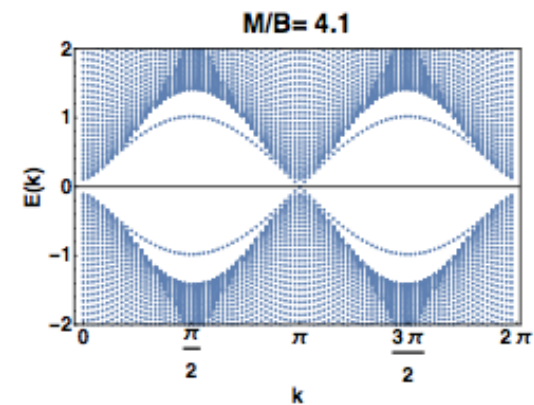
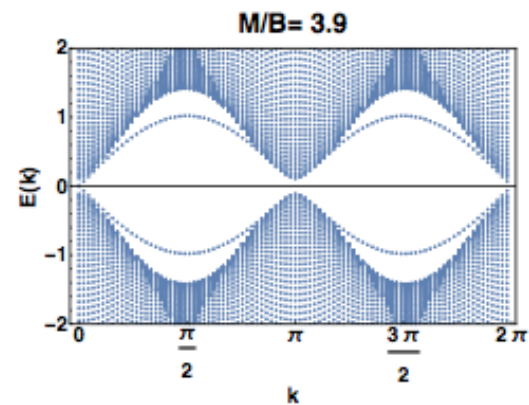
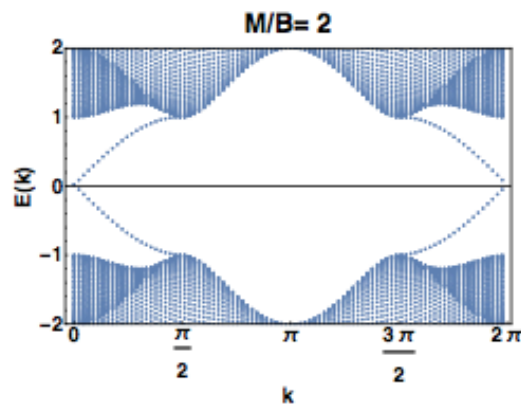
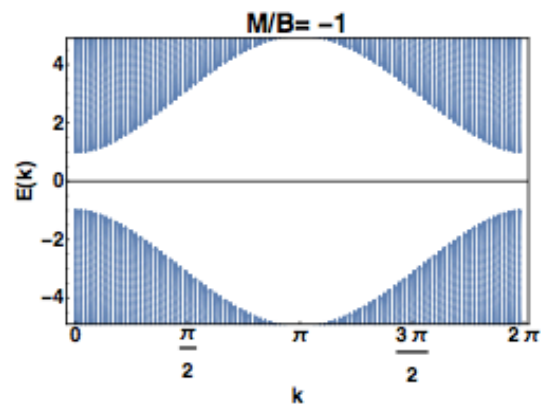
$$0 < M/B < 4d, \quad d = 8$$



$$V_0/B = -6 \times 1$$

Tight Binding (60 x 90 array) Line Impurity

$$0 < M/B < 4d. = 8$$



$$V_0/B = -100 \times 1$$

Conclusions

- Impurities induce local signature of topological insulator
- Experimental signature: probe by applying a gate voltage
 - Normal insulator has an in-gap state for specific voltages
 - TI's will **always** have in-gap state